

Problem Assignment # 6

11/01/2023
due 11/08/20231.4.9. Lorentz transformations in M_2

Consider the 2-dimensional Minkowski space M_2 with metric $g_{ij} = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$ and 2×2 matrix representations of the pseudo-orthogonal group $O(1, 1)$ that leaves g invariant.

a) Let $\sigma, \tau = \pm 1$, and $\phi \in \mathbb{R}$. Show that any element of $O(1, 1)$ can be written in the form

$$D_{\sigma, \tau}(\phi) = \begin{pmatrix} 1 & 0 \\ 0 & \tau \end{pmatrix} \begin{pmatrix} \cosh \phi & \sinh \phi \\ \sinh \phi & \cosh \phi \end{pmatrix} \begin{pmatrix} \sigma & 0 \\ 0 & 1 \end{pmatrix}$$

To study $O(1, 1)$ it thus suffices to study the matrices $D(\phi) := D_{+1, +1} = \begin{pmatrix} \cosh \phi & \sinh \phi \\ \sinh \phi & \cosh \phi \end{pmatrix}$.

b) Show explicitly that the set $\{D(\phi)\}$ forms a group under matrix multiplication (which is a subgroup of $O(1, 1)$ that is sometimes denoted by $SO^+(1, 1)$), and that the mapping $\phi \rightarrow D(\phi)$ defines an isomorphism between this group and the group of real numbers under addition.

c) Show that there exists a matrix J (called the *generator* of the subgroup) such that every $D(\phi)$ can be written in the form

$$D(\phi) = e^{J\phi}$$

and determine J explicitly.

(6 points)

1.4.11. Special Lorentz transformations in M_4

Consider the Minkowski space M_4 .

a) Show that the following transformations are Lorentz transformations:

i) $D^\mu_\nu = \begin{pmatrix} 1 & 0 \\ 0 & R^i_j \end{pmatrix} \equiv R^\mu_\nu$ (rotations)

where R^i_j is any Euclidian orthogonal transformation.

ii) $D^\mu_\nu = \begin{pmatrix} \cosh \alpha & \sinh \alpha & 0 & 0 \\ \sinh \alpha & \cosh \alpha & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \equiv B^\mu_\nu$ (Lorentz boost along the x -direction)

with $\alpha \in \mathbb{R}$.

iii) $D^\mu_\nu = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} \equiv P^\mu_\nu$ (parity)

iv) $D^\mu_\nu = \begin{pmatrix} -1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix} \equiv T^\mu_\nu$ (time reversal)

... /over

- b) Let L be the group of all Lorentz transformations. Show that the rotations defined in part a) i) are a subgroup of L , and so are the Lorentz boosts defined in part a) ii).
- c) Let $I^\mu_\nu = \delta^\mu_\nu$ be the identity transformation. Show that the sets $\{I, P\}$, $\{I, T\}$, and $\{I, P, T, PT\}$ are subgroups of L .

(4 points)

1.5.1. Transformations of tensor fields

- a) Consider a covariant rank- n tensor field $t_{i_1 \dots i_n}(x)$ and find its transformation law under normal coordinate transformations that is analogous to §5.1 def.1; i.e., find how $\tilde{t}_{i_1 \dots i_n}(\tilde{x})$ is related to $t_{i_1 \dots i_n}(x)$.
- b) Convince yourself that your result is consistent with the transformation properties of (i) a covector x_i (the case $n = 1$), and (ii) the covariant components of the metric tensor g_{ij} .

(4 points)

1.5.2. Curl and divergence

Show that the curl and the divergence of a vector field transform as a pseudovector field and a scalar field, respectively.

(3 points)

1.5.3. Tensor products, and tensor traces

Prove Propositions 1 and 2 from ch. 1 §5.3.

(4 points)

1.4.9, a) Let $D = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$. For D to be a Lorentz transformation, we must have

$$\begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} = \begin{pmatrix} a & b \\ c & d \end{pmatrix} \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \begin{pmatrix} a & c \\ b & d \end{pmatrix} = \begin{pmatrix} a^2 - b^2 & ac - bd \\ ac - bd & c^2 - d^2 \end{pmatrix}$$

From this obtain three constraints on the four numbers a, b, c, d

$$(i) \quad a^2 - b^2 = 1$$

$$(ii) \quad c^2 - d^2 = -1$$

$$(iii) \quad ac - bd = 0$$

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Now consider $\text{nil } \phi$, which maps \mathbb{R} one-to-one into itself

$$\rightarrow \forall b \in \mathbb{R} \exists! \phi \in \mathbb{R} : \quad \underline{b = \text{nil } \phi}$$

$$(i) \rightarrow a^2 = 1 + b^2 = 1 + \text{nil}^2 \phi = \text{nil}^2 \phi \rightarrow \underline{a = \tau \text{nil } \phi}, \quad \underline{\tau = \pm 1}$$

Analogously, $\underline{c = \tau' \text{nil } \phi}$

$$(ii) \rightarrow d^2 = 1 + c^2 = 1 + \text{nil}^2 \phi = \text{nil}^2 \phi \rightarrow \underline{d = \tau' \text{nil } \phi}, \quad \underline{\tau' = \pm 1}$$

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Finally,

$$(iii) \rightarrow 0 = \tau \text{nil } \phi \text{ nil } \phi - \tau' \text{nil } \phi \text{ nil } \phi$$

$$= \text{nil } \phi \text{ nil } \phi - \tau \tau' \text{nil } \phi \text{ nil } \phi$$

$$= \text{nil} (\tau \tau' \phi) \text{ nil } \phi - \text{nil } \phi \text{ nil} (\tau \tau' \phi)$$

$$= \text{nil} (\phi - \tau \tau' \phi) \rightarrow \underline{\phi = \tau \tau' \phi}$$

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$$\rightarrow \underline{\Delta_{\tau, \tau'}(\phi)} = \begin{pmatrix} \tau \text{nil } \phi & \text{nil } \phi \\ \text{nil}(\tau \tau' \phi) & \tau \text{nil}(\tau \tau' \phi) \end{pmatrix} = \begin{pmatrix} \tau \text{nil } \phi & \text{nil } \phi \\ \tau \tau' \text{nil } \phi & \tau \text{nil } \phi \end{pmatrix}$$

$$= \begin{pmatrix} 1 & 0 \\ 0 & \tau \end{pmatrix} \begin{pmatrix} \tau \text{nil } \phi & \text{nil } \phi \\ \tau \text{nil } \phi & \text{nil } \phi \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & \tau \end{pmatrix} \begin{pmatrix} \text{nil } \phi & \text{nil } \phi \\ \text{nil } \phi & \text{nil } \phi \end{pmatrix} \begin{pmatrix} \tau & 0 \\ 0 & 1 \end{pmatrix}$$

$$= \begin{pmatrix} 1 & 0 \\ 0 & \tau \end{pmatrix} \underline{\Delta(\phi)} \begin{pmatrix} \tau & 0 \\ 0 & 1 \end{pmatrix} \quad \text{with} \quad \underline{\Delta(\phi)} = \begin{pmatrix} \text{nil } \phi & \text{nil } \phi \\ \text{nil } \phi & \text{nil } \phi \end{pmatrix}, \quad \phi \in \mathbb{R}$$

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and $\tau, \tau' = \pm 1$ is the most general element of $O(1,1)$.

$$b) \quad (i) \quad \begin{pmatrix} \cos \phi_1 & \sin \phi_1 \\ \sin \phi_1 & \cos \phi_1 \end{pmatrix} \begin{pmatrix} \cos \phi_2 & \sin \phi_2 \\ \sin \phi_2 & \cos \phi_2 \end{pmatrix} = \begin{pmatrix} \cos(\phi_1 + \phi_2) & \sin(\phi_1 + \phi_2) \\ \sin(\phi_1 + \phi_2) & \cos(\phi_1 + \phi_2) \end{pmatrix}$$

$$\rightarrow \Delta(\phi_1)\Delta(\phi_2) = \Delta(\phi_1 + \phi_2) \quad \text{down } \checkmark$$

(ii) Matrix multiplication is associative \checkmark

$$(iii) \quad \Delta(\phi=0) = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} = \mathbb{1}_2 \quad \text{neutral element } \checkmark$$

$$(iv) \quad \Delta(-\phi) = \begin{pmatrix} \cos \phi & -\sin \phi \\ \sin \phi & \cos \phi \end{pmatrix}$$

$$\text{and } \begin{pmatrix} \cos \phi & -\sin \phi \\ -\sin \phi & \cos \phi \end{pmatrix} \begin{pmatrix} \cos \phi & \sin \phi \\ \sin \phi & \cos \phi \end{pmatrix} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix}$$

$$\rightarrow \Delta(-\phi) = (\Delta(\phi))^{-1} \quad \text{inverse } \checkmark$$

$$\rightarrow \underline{\{\Delta(\phi)\} \text{ is a group } SO^+(1,1)}$$

(i), (iii), (iv) provide the isomorphism $\underline{SO^+(1,1) \cong \mathbb{R}(+)}$

$$c) \quad e^{\gamma \phi} = \mathbb{1}_2 + \gamma \phi + \frac{1}{2} \gamma^2 \phi^2 + \dots$$

$$\text{and } \cos \phi = 1 + \frac{1}{2} \phi^2 + \frac{1}{4!} \phi^4 + \dots, \quad \sin \phi = \phi + \frac{1}{3!} \phi^3 + \dots$$

$$\text{try } \underline{\gamma = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}} \rightarrow \gamma^2 = \mathbb{1}_2, \quad \gamma^3 = \gamma, \text{ etc.}$$

$$\rightarrow \underline{e^{\gamma \phi} = \begin{pmatrix} 1 & 0 \\ 0 & 1 \end{pmatrix} + \begin{pmatrix} 0 & \phi \\ \phi & 0 \end{pmatrix} + \begin{pmatrix} \frac{1}{2} \phi^2 & 0 \\ 0 & \frac{1}{2} \phi^2 \end{pmatrix} + \begin{pmatrix} 0 & \frac{1}{3!} \phi^3 \\ \frac{1}{3!} \phi^3 & 0 \end{pmatrix} + \dots}$$

$$= \underline{\underline{\begin{pmatrix} \cos \phi & \sin \phi \\ \sin \phi & \cos \phi \end{pmatrix}}} \quad \checkmark$$

$\rightarrow \underline{\underline{\gamma}}$ is the generator of $\underline{SO^+(1,1)}$.

1.4.11.2.0) i)
$$\underline{R^T \times g_{\mu\nu} R^\nu}_\Lambda = \left(\begin{array}{c|c} 1 & 0 \\ \hline 0 & R^T \end{array} \right) \left(\begin{array}{c|c} 1 & 0 \\ \hline 0 & -1 \end{array} \right) \left(\begin{array}{c|c} 1 & 0 \\ \hline 0 & R \end{array} \right) = \left(\begin{array}{c|c} 1 & 0 \\ \hline 0 & R^T \end{array} \right) \left(\begin{array}{c|c} 1 & 0 \\ \hline 0 & -R \end{array} \right)$$

$$= \left(\begin{array}{c|c} 1 & 0 \\ \hline 0 & -R^T R \end{array} \right) = \left(\begin{array}{c|c} 1 & 0 \\ \hline 0 & -1 \end{array} \right) = \underline{g_{\Lambda\Lambda}} \quad \checkmark$$

ii)
$$\underline{\Pi^T \times g_{\mu\nu} \Pi^\nu}_\Lambda = \begin{pmatrix} w_1 L_1 & 0 & 0 & w_1 L_1 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ w_1 L_1 & 0 & 0 & w_1 L_1 \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix}$$

$$\times \begin{pmatrix} w_1 L_1 & 0 & 0 & w_1 L_1 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ w_1 L_1 & 0 & 0 & w_1 L_1 \end{pmatrix}$$

$$= \begin{pmatrix} w_1^2 L_1^2 - w_1^2 L_1^2 & 0 & 0 & w_1^2 L_1^2 - w_1^2 L_1^2 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ w_1^2 L_1^2 - w_1^2 L_1^2 & 0 & 0 & w_1^2 L_1^2 - w_1^2 L_1^2 \end{pmatrix}$$

$$= \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix} = \underline{g_{\Lambda\Lambda}} \quad \checkmark$$

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iii)
$$\underline{P^T \times g_{\mu\nu} P^\nu}_\Lambda = \left(\begin{array}{c|c} 1 & 0 \\ \hline 0 & -1 \end{array} \right) \left(\begin{array}{c|c} 1 & 0 \\ \hline 0 & -1 \end{array} \right) \left(\begin{array}{c|c} 1 & 0 \\ \hline 0 & -1 \end{array} \right) = \left(\begin{array}{c|c} 1 & 0 \\ \hline 0 & -1 \end{array} \right) = \underline{g_{\Lambda\Lambda}} \quad \checkmark$$

iv)
$$\underline{T^T \times g_{\mu\nu} T^\nu}_\Lambda = (-1)^2 P^T \times g_{\mu\nu} P^\nu_\Lambda = \underline{g_{\Lambda\Lambda}} \quad \checkmark$$

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b) R^T_ν leaves the time coordinates invariant, and we know that the R^i_j form a group (in $PL(3,1) \cong O(3,1)$) $\rightarrow \{R^i_j\}$ is a subgroup of L .
 The Lorentz boosts \mathcal{B}_ν on a subgroup of L according to Problem 1.8.

c) group tables: $\begin{array}{c|c} \mathcal{B} & \mathcal{B} \\ \hline \mathcal{B} & \mathcal{B} \\ \mathcal{P} & \mathcal{P} \end{array}$ and $\begin{array}{c|c} \mathcal{B} & \mathcal{T} \\ \hline \mathcal{B} & \mathcal{T} \\ \mathcal{T} & \mathcal{T} \end{array} \rightarrow \{\mathcal{B}, \mathcal{P}\}$, and $\{\mathcal{B}, \mathcal{T}\}$ are subgroups of L with matrix multiplication

	\mathcal{B}	\mathcal{P}	\mathcal{T}	\mathcal{PT}
\mathcal{B}	\mathcal{B}	\mathcal{P}	\mathcal{T}	\mathcal{PT}
\mathcal{P}	\mathcal{P}	\mathcal{B}	\mathcal{PT}	\mathcal{T}
\mathcal{T}	\mathcal{T}	\mathcal{PT}	\mathcal{B}	\mathcal{P}
\mathcal{PT}	\mathcal{PT}	\mathcal{T}	\mathcal{P}	\mathcal{B}

$\rightarrow \{\mathcal{B}, \mathcal{P}, \mathcal{T}, \mathcal{PT}\}$ is also a subgroup of L

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1.5.1. a) There are various ways to do this. One option is to start with the transformation property of covariant tensor fields, §5.1, and use the metric tensor to lower the indices:

$$\begin{aligned}
 \underline{\tilde{T}}_{i_1 \dots i_N}(\bar{x}) &= \tilde{g}_{i_1 j_1} \dots \tilde{g}_{i_N j_N} \tilde{T}^{j_1 \dots j_N}(\bar{x}) \\
 &\stackrel{\text{§5.1(4)} + \text{§5.3}}{=} \tilde{g}_{i_1 j_1} \dots \tilde{g}_{i_N j_N} \Delta^{j_1}_{k_1} \dots \Delta^{j_N}_{k_N} T^{k_1 \dots k_N}(x) \\
 &= \tilde{g}_{i_1 j_1} \dots \tilde{g}_{i_N j_N} \Delta^{j_1}_{k_1} \dots \Delta^{j_N}_{k_N} g^{k_1 l_1} \dots g^{k_N l_N} T_{l_1 \dots l_N}(x) \\
 &= (g \Delta)_{i_1 k_1} \dots (g \Delta)_{i_N k_N} g^{k_1 l_1} \dots g^{k_N l_N} T_{l_1 \dots l_N}(x) \\
 &\stackrel{g \Delta = \Delta^T g}{=} ((\Delta^T)^{-1})_{i_1 k_1} \dots ((\Delta^T)^{-1})_{i_N k_N} g^{k_1 l_1} \dots g^{k_N l_N} T_{l_1 \dots l_N}(x) \\
 &= ((\Delta^T)^{-1})_{i_1}^{m_1} \dots ((\Delta^T)^{-1})_{i_N}^{m_N} g^{m_1 l_1} \dots g^{m_N l_N} T_{l_1 \dots l_N}(x) \\
 &= ((\Delta^T)^{-1})_{i_1}^{m_1} \underbrace{g^{m_1 l_1}}_{\delta_{m_1 l_1}} \dots ((\Delta^T)^{-1})_{i_N}^{m_N} \underbrace{g^{m_N l_N}}_{\delta_{m_N l_N}} T_{l_1 \dots l_N}(x) \\
 &= ((\Delta^T)^{-1})_{i_1}^{j_1} \dots ((\Delta^T)^{-1})_{i_N}^{j_N} T_{j_1 \dots j_N}(x)
 \end{aligned}$$

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This is §5.1 (2) with $\Delta^{i_n}_{j_n}$ replaced by $((\Delta^T)^{-1})_{i_n}^{j_n}$ ($n=1, \dots, N$)

b) Special case $n=1$:

$$\begin{aligned}
 \underline{\tilde{x}}_i &= \tilde{g}_{ij} \tilde{x}^j \stackrel{\text{§4.8.2}}{=} \tilde{g}_{ij} \Delta^j_k x^k \stackrel{\tilde{g}^{-1}}{=} (g \Delta)_{ik} x^k \stackrel{g \Delta = \Delta^T g}{=} ((\Delta^T)^{-1})_{ik} x^k \\
 &= ((\Delta^T)^{-1})_{i j} \tilde{x}^j \quad \underline{((\Delta^T)^{-1})_{i j} x^j} \quad \checkmark
 \end{aligned}$$

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$$\begin{aligned}
 \tilde{g}_{ij} &= g_{ik} g_{jl} \tilde{g}^{kl} \stackrel{\text{§5.1}}{=} g_{ik} g_{jl} \Delta^k_m \Delta^l_n g^{mn} = (g \Delta)_{im} (\Delta^T)_jn g^{mn} \\
 &= ((\Delta^T)^{-1})_{i m} ((\Delta^T)^{-1})_{j n} g^{mn}
 \end{aligned}$$

1.5.2.) Consider the curl as defined in § 5.2 :

$$c^i(x) = \epsilon^{ijkl} \partial_j v_k(x)$$

$$\rightarrow \tilde{c}^i(\tilde{x}) = \tilde{\epsilon}^{ijkl} \tilde{\partial}_j \tilde{v}_k(\tilde{x})$$

Proof 1.5.2

$$= \tilde{\epsilon}^{ijkl} (\Delta^{-1})^e_j \partial_e (\Delta^{-1})^m_k v_m(x)$$

$$= \delta^i_n \tilde{\epsilon}^{nijk} (\Delta^{-1})^e_j (\Delta^{-1})^m_k \partial_e v_m(x)$$

$$= \Delta^i_p (\Delta^{-1})^p_n \tilde{\epsilon}^{nijk} (\Delta^{-1})^e_j (\Delta^{-1})^m_k \partial_e v_m(x)$$

$$= \Delta^i_p (\Delta^{-1})^p_n (\Delta^{-1})^e_j (\Delta^{-1})^m_k \tilde{\epsilon}^{nijk} \partial_e v_m(x)$$

$$= (\det \Delta) \epsilon^{pkmn} \text{ by § 5.1 remark (9)}$$

$$= (\det \Delta) \Delta^i_p \underbrace{\epsilon^{pkmn} \partial_e v_m(x)}_{= c^p(x)}$$

$$= (\det \Delta) \Delta^i_p c^p(x)$$

$\tilde{c}^i(x)$ transforms as a pseudovector field.

Now the divergence: $d(x) = \partial_i v^i(x)$

$$\rightarrow \tilde{d}(\tilde{x}) = \tilde{\partial}_i \tilde{v}^i(\tilde{x}) = ((\Delta^{-1})^{-1})^i_j \partial_j \Delta^i_k v^k(x)$$

$$= (\Delta^T)_k^i ((\Delta^T)^{-1})^j_i \partial_j v^k(x)$$

$$= \delta_k^j$$

$$= \partial_k v^k(x) \rightarrow \underline{\underline{d(x) transforms as a scalar field}}$$

1.5.1) Consider the lower product

$$u^{i_1 \dots i_N} = s^{i_1 \dots i_N} + t^{i_1 \dots i_N}$$

Let both s and t be lower. Then

$$\begin{aligned} \underline{\tilde{u}^{i_1 \dots i_N}} &= \tilde{s}^{i_1 \dots i_N} + \tilde{t}^{i_1 \dots i_N} \\ &= \prod_{j_1}^{i_1} \dots \prod_{j_N}^{i_N} s^{j_1 \dots j_N} + \prod_{j_1}^{i_1} \dots \prod_{j_N}^{i_N} t^{j_1 \dots j_N} \\ &= \prod_{j_1}^{i_1} \dots \prod_{j_N}^{i_N} u^{j_1 \dots j_N} \quad (*) \end{aligned}$$

$\Rightarrow u$ is a rank- (N, N) lower

Now suppose s is a pseudo lower and t a lower, or vice versa. Then the rhs of (*) gets multiplied by $\det D \Rightarrow$ u is a pseudo lower

If both s and t are pseudo lowers, then the rhs of (*) gets multiplied by $(\det D)^2 = I \Rightarrow$ u is a lower

This proves prop. 1 from § 5.2.

Now consider

$$u^{k_1 \dots k_N} = t^{i_1 k_1 \dots k_N} = \prod_{j_1}^{i_1} t^{j_1 k_1 \dots k_N}$$

Let t be a lower. Then

$$\underline{\tilde{u}^{k_1 \dots k_N}} = \tilde{\prod}_{j_1}^{i_1} t^{j_1 k_1 \dots k_N}$$

Proof: $\tilde{u}^{k_1 \dots k_N} = ((D^{-1})^T)^0_i ((D^{-1})^T)^P_j \prod_{op} D^i_m D^j_n D^{k_1} \dots D^{k_N} t^{m k_1 \dots k_N}$
 $= \underbrace{(D^{-1})^0_i D^i_m}_{= \delta^0_m} \underbrace{(D^{-1})^P_j D^j_n}_{= \delta^P_n} D^{k_1} \dots D^{k_N} \prod_{op} t^{m k_1 \dots k_N}$

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p1.5.3-2

$$= \Delta_{l_1}^{h_1} \dots \Delta_{l_N}^{h_N} \sum_{u_{l_1 \dots l_N}} t$$

$$= \Delta_{l_1}^{h_1} \dots \Delta_{l_N}^{h_N} u_{l_1 \dots l_N} \quad (10)$$

\Rightarrow u is a lower of rank N

$\exists!$ t is a pseudoscalar, then the r.h.s of (10) gets multiplied by $\det \Delta$

\Rightarrow u is a pseudoscalar of rank N

1
This proves prop 2 for §5.3