

0.3.1. Particle in homogeneous $\mathbf{E}$ and $\mathbf{B}$ fields

Consider a point particle (mass m , charge e) in homogeneous fields $\mathbf{B} = (0, 0, B)$ and $\mathbf{E} = (0, E_y, E_z)$. Treat the motion of the particle nonrelativistically.

a) Show that the motion in z -direction decouples from the motion in the x - y plane, and find $z(t)$.

b) Consider $\xi := x + iy$. Find the equation of motion for ξ , and its most general solution.

hint: Define the *cyclotron frequency* $\omega = eB/mc$, and remember how to solve inhomogeneous ODEs.

c) Show that the time-averaged velocity perpendicular to the plane defined by $\mathbf{B}$ and $\mathbf{E}$ is given by the *drift velocity*

$$\langle \mathbf{v} \rangle = c \mathbf{E} \times \mathbf{B} / B^2$$

Show that $E_y/B \ll 1$ is necessary and sufficient for the non relativistic approximation to be valid.

d) Show that the path projected onto the x - y plane can have three qualitatively different shapes, and plot a representative example for each.

(6 points)

0.3.1.) a) Eq. of motion:

$$m\vec{\dot{v}} = e\vec{E} + \frac{e}{c}\vec{v} \times \vec{B}$$

$$\text{let } \vec{B} = (0, 0, B) \text{ and } \vec{E} = (0, E_y, E_z)$$

$$\rightarrow \begin{cases} m\ddot{x} = \frac{e}{c} \dot{y} B & (1) \\ m\ddot{y} = eE_y - \frac{e}{c} \dot{x} B & (2) \\ m\ddot{z} = eE_z & (3) \end{cases}$$

$$(3) \rightarrow \underline{\underline{z(t) = z_0 + v_z^0 t + \frac{eE_z}{2m} t^2}}$$

b) Define $\underline{\underline{\zeta := x + iy}}$

$$(1) + i \cdot (2) \rightarrow m\ddot{\zeta} = ieE_y - i\frac{eB}{c}\dot{\zeta}$$

Define $\omega := \frac{eB}{mc}$ cyclotron frequency

$$\rightarrow \boxed{\ddot{\zeta} + i\omega\dot{\zeta} = i\frac{e}{m}E_y} \quad (*)$$

Special solution of inhomogeneous eq: $\dot{\zeta} = \frac{eE_y}{m\omega}$ General solution of homogeneous eq: $\dot{\zeta} = ae^{-i\omega t} \quad (a \in \mathbb{C})$

$\rightarrow \underline{\underline{\dot{\zeta}(t) = ae^{-i\omega t} + eE_y/m\omega}}$ is the most general solution of (*).

c) With $a = be^{i\alpha}$, $b, \alpha \in \mathbb{R}$

$$\rightarrow \dot{\zeta} = be^{-i(\omega - \alpha)t} + eE_y/m\omega$$

$\rightarrow \alpha$ just shifts the start of time $\rightarrow \underline{\underline{\alpha = 0 \text{ w.l.o.g.}}}$

$$\rightarrow \dot{x} + i\dot{y} = b\omega e^{i\omega t} - ib\omega e^{i\omega t} + eE_y/m\omega$$

$$\rightarrow \begin{cases} \dot{x} = b\omega e^{i\omega t} + eE_y/m\omega \\ \dot{y} = -b\omega e^{i\omega t} \end{cases} \quad (**)$$

①

$$\rightarrow \langle \dot{y} \rangle = 0, \quad \langle \dot{x} \rangle = eE_y/m\omega = \frac{eE_y}{\Omega} \quad \text{time-averaged velocity}$$

$$= \frac{eE_y \Omega}{\Omega^2} = \frac{c(\vec{E} \times \vec{\Omega})_x}{\Omega^2}$$

in general:

$$\langle \vec{v} \rangle = \frac{c}{\Omega^2} \vec{E} \times \vec{\Omega} \quad \text{drift velocity}$$

①

$$\text{condition for } v \ll c: \quad \underline{E_y/\Omega \ll 1}$$

necessary and sufficient
condition for non-
relativistic approximation

$$d) \quad \text{Given } x(t=0) = 0 = y(t=0) \quad \text{w.l.o.g.}$$

$$(**) \rightarrow \begin{cases} x(t) = \frac{b}{\omega} \omega e^{i\omega t} + \frac{eE_y}{\Omega} t \\ y(t) = \frac{b}{\omega} (\omega e^{i\omega t} - 1) \end{cases}$$

$\rightarrow$ The path is a trochoid.

To visualize it, put $\omega = 1$ and define $C = eE_y/\Omega$.

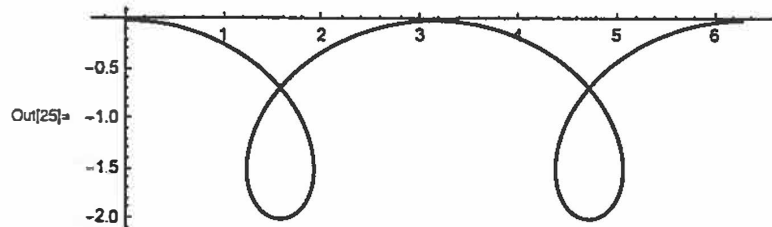
$$\rightarrow \begin{cases} x(t) = b e^{it} + Ct \\ y(t) = b (e^{it} - 1) \end{cases}$$

This is the projection of
the path onto the x-y plane

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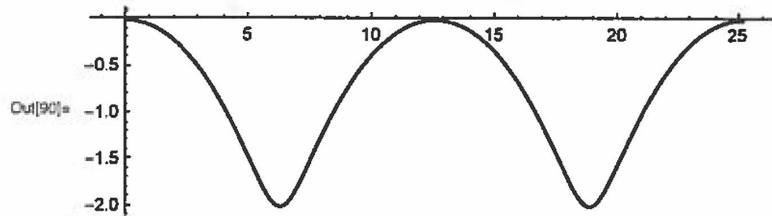
For $C < b$ the trochoid has loops :

```
In[21] = b = 1;
c = 0.5;
x[t_] := b Sin[t] + c t
y[t_] := b (Cos[t] - 1)
ParametricPlot[{x[t], y[t]}, {t, 0, 4 Pi}]
```



For $C > b$ it does not :

```
In[88] = b = 1;
c = 2;
x[t_] := b Sin[t] + c t
y[t_] := b (Cos[t] - 1)
ParametricPlot[{x[t], y[t]}, {t, 0, 4 Pi}, AspectRatio -> 0.3]
```



And for $C = b$ it degenerates into a cycloid :

```
In[91] = b = 1;
c = 1;
x[t_] := b Sin[t] + c t
y[t_] := b (Cos[t] - 1)
ParametricPlot[{x[t], y[t]}, {t, 0, 4 Pi}, AspectRatio -> 0.3]
```

