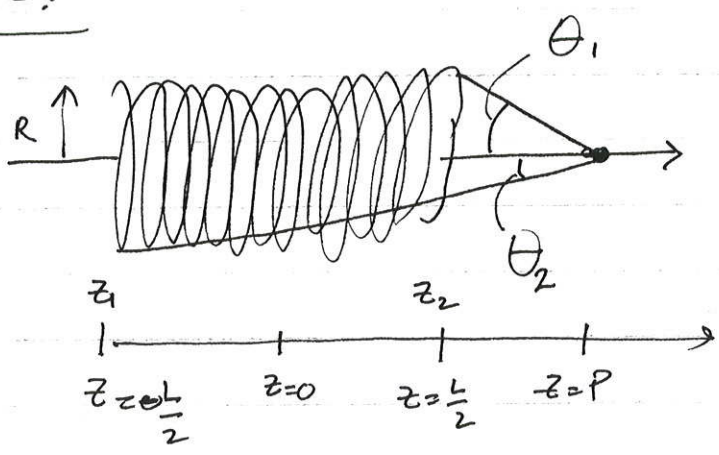


# HW #2

## Prob 5.11



Find the  $\vec{B}$ -field on the axis of a tightly wound solenoid consisting of  $N$  turns per unit length wrapped around a cylindrical tube of radius  $R$  and carrying current  $I$ .

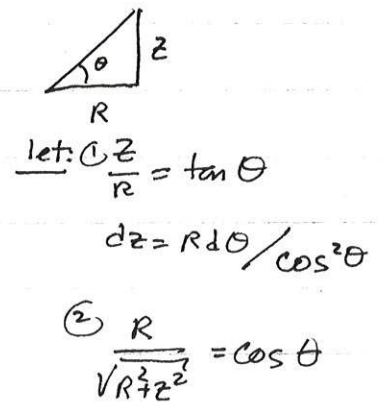
Example 6  $\Rightarrow B_z = \frac{\mu_0 I}{2} \frac{R^2}{(R^2 + z^2)^{3/2}}$  (5.34)

for the axis of a circular loop.

$\vec{J} =$  current per unit length along  $z = NI = \left[ \frac{A}{m} \right]$

$$d\vec{B}_s = \frac{\mu_0 \vec{J} dz R^2}{2 (R^2 + z^2)^{3/2}} \hat{z}$$

$$\begin{aligned} \vec{B}_s &= \frac{\mu_0 N I R^2}{2} \int_{z_1}^{z_2} \frac{dz}{(R^2 + z^2)^{3/2}} \\ &= \frac{\mu_0 N I R^2}{2} \int_{R \tan \theta_1}^{R \tan \theta_2} \frac{R d\theta \cos^3 \theta}{\cos^2 \theta R^3} \\ &= \frac{\mu_0 N I}{2} (\sin \theta) \Big|_{R \tan \theta_1}^{R \tan \theta_2} \end{aligned}$$



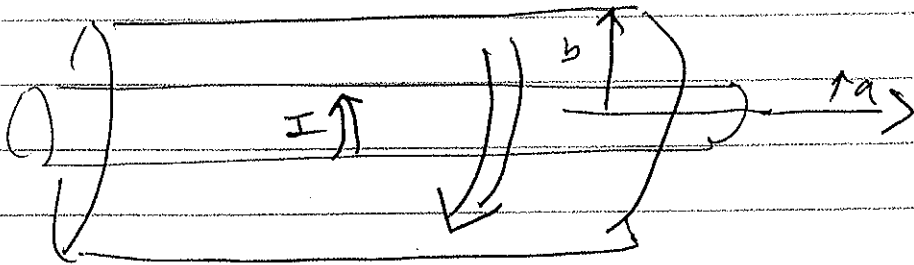
$$\begin{aligned} &= \frac{\mu_0 N I}{2} \left[ \frac{-z_1}{\sqrt{R^2 + z_1^2}} + \frac{z_2}{\sqrt{R^2 + z_2^2}} \right] \\ &= \frac{\mu_0 N I}{2} \left[ \frac{L/2}{\sqrt{R^2 + L^2/4}} - \frac{-L/2}{\sqrt{R^2 + L^2/4}} \right] \end{aligned}$$

$$\vec{B}_s = \frac{\mu_0 N I L}{2 \sqrt{R^2 + L^2/4}} \hat{z}$$

if  $L \rightarrow \infty$ ,

$$\vec{B}_s \approx \mu_0 N I \hat{z}$$

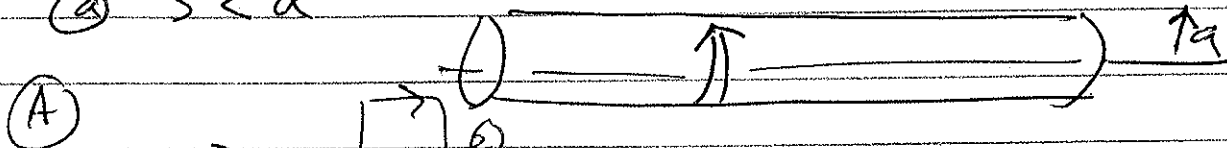
Prob 5.15



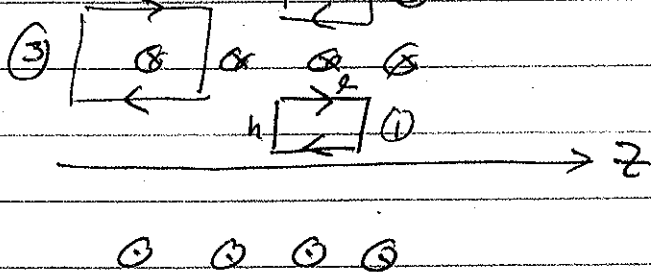
find  $\vec{B}$  in (i)  $s < a$ , (ii)  $a < s < b$ , (iii)  $b < s$

Sol<sup>n</sup>

(a)  $s < a$



(A)



$$\oint \vec{B} \cdot d\vec{\ell} = \mu_0 \int \vec{J} \cdot d\vec{s}$$

$$\textcircled{1} B_{z,<}^+ l - B_{z,<}^b l = 0$$

$\Rightarrow B_{z,<}^+ \text{ is constant}$

$$\textcircled{2} B_{z,>}^+ l - B_{z,>}^b l = 0$$

as  $B_{z,>} \rightarrow 0$  at  $\infty$

$$\Rightarrow B_{z,>} = 0$$

$$\textcircled{3} B_{z,>}^+ l - B_{z,<}^b l = N \mu_0 I$$

$$\boxed{\vec{B}_{z,<} = -\mu_0 I \hat{z}}$$

(B) Outer cylinder contributes,  $\vec{B}_{z,>}^{\text{outer}} = +\mu_0 I \hat{z}$

$$\Rightarrow \vec{B}_{\text{total}} = \mu_0 I (n_2 - n_1) \hat{z}, \quad s < a$$

(B) Between,  $a < s < b$

$$\vec{B}_{\text{total}} = \vec{B}_{\text{inner}} + \vec{B}_{\text{outer}}$$

$$= 0 + n_2 \mu_0 I \hat{z}$$

$$\vec{B}_{\text{total}} = \mu_0 n_2 I \hat{z}, \quad a < s < b$$

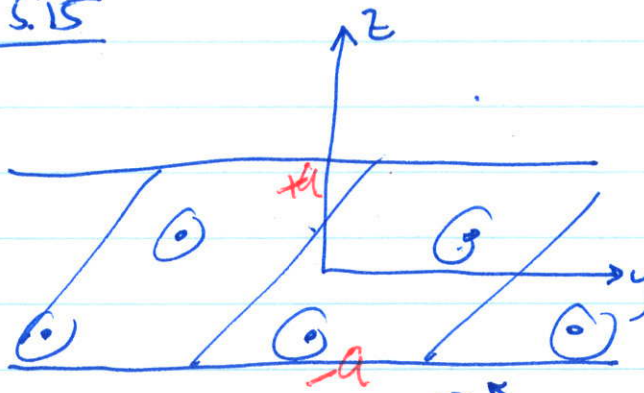
(C) Outside,  $b < s$

$$\vec{B}_{\text{total}} = \vec{B}_{\text{inner}} + \vec{B}_{\text{outer}}$$

$$= 0 + 0$$

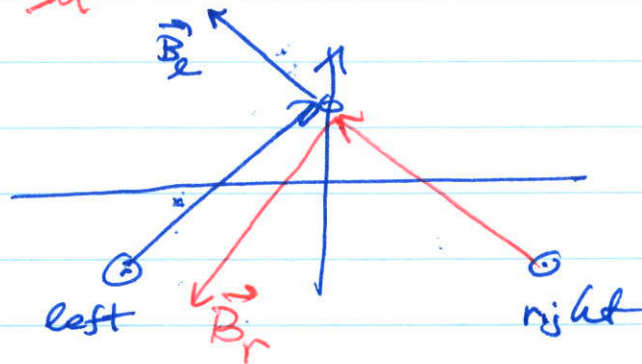
$$\vec{B}_{\text{total}} = 0, \quad b < s$$

Prob 5.15



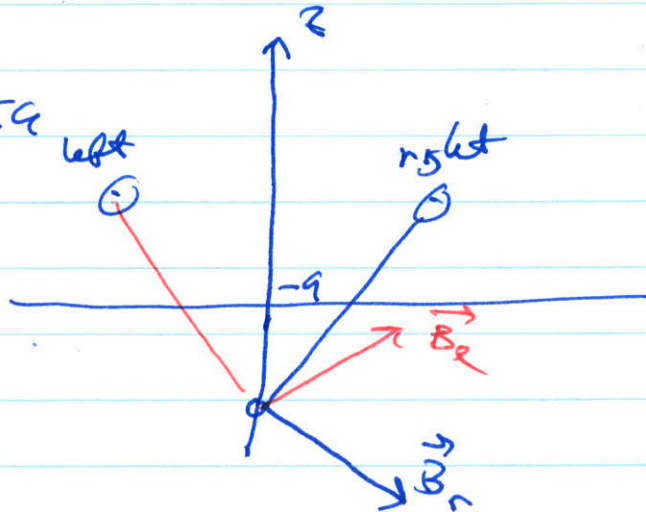
$\vec{J} = J_0 \hat{x}$   
for  $\vec{B}$  everywhere

a) for  $z > a$



$\Rightarrow$  only  $B_y$  survives,  $B_z$  cancels. also  $B_x = 0$  b/c  $\vec{J}$  is in x direction

b) for  $z < a$

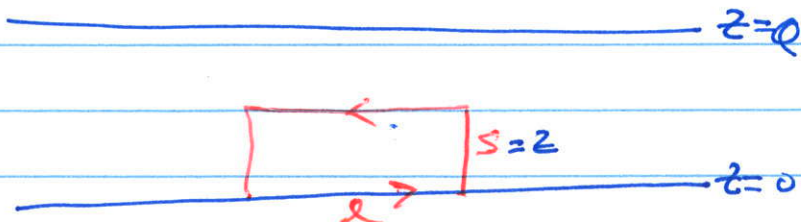


$\Rightarrow$  only  $B_y$  survives but is in opposite direction compared to the  $z > a$  field.

c)  $\Rightarrow$  at  $z=0$ ,  $B_y=0$  (by symmetry)

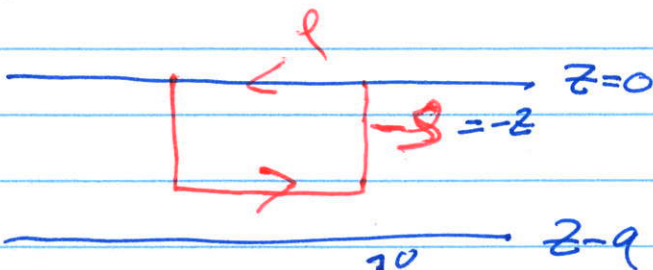
$$\oint \vec{B} \cdot d\vec{l} = \mu_0 \int \vec{J} \cdot d\vec{S}$$

for  $0 < z < a$



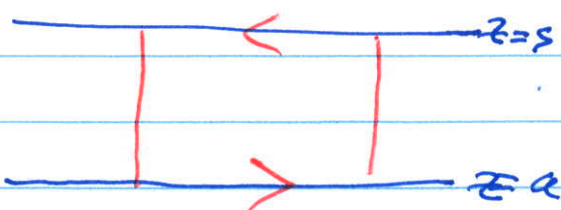
$$-B_y(s)l + B_y(0)l = \mu_0 J_0 l s \Rightarrow B_y(s) = -\mu_0 J_0 s$$

for  $-a < z < 0$



$$+B_y(-s)l + B_y(0)(-l) = +\mu_0 J_0 l s \Rightarrow B_y(s) = \mu_0 J_0 s$$

for  $a < z < s$



$$\Rightarrow -B_y(s)l + B_y(a)l = 0 \Rightarrow B_y(s) = B_y(a)$$

for  $-a > z > -s$

$$\Rightarrow B_y(-s) = B_y(-a)$$



write this in a more sensible way using  $z$ , then

$$B_y(z) = \begin{cases} -\mu_0 J_0 a & z > a \\ -\mu_0 J_0 z & -a < z < a \\ \mu_0 J_0 a & z < -a \end{cases}$$

## Prob 5.22

Suppose monopoles existed. How would you modify Maxwell's equations and the force law?

Current Maxwell & Lorentz law

$$\begin{aligned}\vec{\nabla} \cdot \vec{E} &= \rho/\epsilon_0 & \vec{\nabla} \cdot \vec{B} &= 0 \\ \vec{\nabla} \times \vec{E} &= 0 & \vec{\nabla} \times \vec{B} &= \mu_0 \vec{J}\end{aligned} \quad \& \quad \vec{F} = q\vec{E} + q\vec{v} \times \vec{B}$$

electrostatics & magnetostatics

a) Maxwell's equations

$$\begin{aligned}\vec{\nabla} \cdot \vec{E} &= \rho/\epsilon_0 & \vec{\nabla} \cdot \vec{B} &= C_1 \rho_m \\ \vec{\nabla} \times \vec{E} &= -C_2 \vec{J}_m & \vec{\nabla} \times \vec{B} &= \mu_0 \vec{J}\end{aligned}$$

$C_1$  and  $C_2$  are not clear (yet). For example, later we write

$$\vec{\nabla} \times \vec{B} = \mu_0 \left[ \vec{J} + \epsilon_0 \frac{\partial \vec{E}}{\partial t} \right]. \quad \text{Then } \vec{\nabla} \cdot \left[ \vec{\nabla} \times \vec{B} \right] = \mu_0 \left[ \vec{\nabla} \cdot \vec{J} + \epsilon_0 \vec{\nabla} \cdot \frac{\partial \vec{E}}{\partial t} \right]$$

$$0 = \mu_0 \vec{\nabla} \cdot \vec{J} + \mu_0 \epsilon_0 \frac{\partial}{\partial t} \vec{\nabla} \cdot \vec{E}$$

$$= \vec{\nabla} \cdot \vec{J} + \frac{\partial \rho}{\partial t} \quad \text{and charge is conserved}$$

$$\text{but } \vec{\nabla} \times \vec{E} = -\frac{\partial \vec{B}}{\partial t} \Rightarrow \vec{\nabla} \times \vec{E} = -C_2 \vec{J}_m - \frac{\partial \vec{B}}{\partial t} \quad \text{w/ monopoles}$$

$$\vec{\nabla} \cdot (\vec{\nabla} \times \vec{E}) = C_2 \vec{\nabla} \cdot \vec{J}_m - \frac{\partial}{\partial t} \vec{\nabla} \cdot \vec{B}$$

$$= C_2 \vec{\nabla} \cdot \vec{J}_m - \frac{\partial}{\partial t} C_1 \rho_m$$

$$= \frac{C_2}{C_1} \left[ \vec{\nabla} \cdot \vec{J}_m \right] - \frac{\partial \rho_m}{\partial t} = 0$$



this will be true if  $\frac{c_2}{c_1} = -1$

(b) force law

$$\vec{F} = (\rho_e \vec{E} + \vec{J}_e \times \vec{B}) + (\rho_m \vec{B} + c_3 \vec{J}_m \times \vec{E})$$

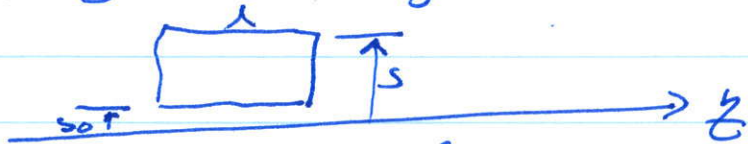
where  $c_3$  is determined by units and charge convention

Prob 5.26

(a) Find  $\vec{A}$  for an infinite line current.

(i)  $\vec{B} = \frac{\mu_0 I}{2\pi s} \hat{\phi}$  if  $z$  and current  $I$  are aligned

(ii)  $\int \vec{B} \cdot d\vec{S} = \int \vec{A} \cdot d\vec{\ell}$



(i)  $\int \vec{B} \cdot d\vec{S} = \frac{\mu_0 I}{2\pi} \int_{s_0}^s dz ds \frac{1}{s} = \frac{\mu_0 I l}{2\pi} \ln\left(\frac{s}{s_0}\right)$

(ii)  $\int \vec{A} \cdot d\vec{\ell} = A_z(s_0)l - A_z(s)l$

$\Rightarrow A_z(s) = A_z(s_0) - \frac{\mu_0 I}{2\pi} \ln\left(\frac{s}{s_0}\right)$

(iii) find  $\vec{\nabla} \cdot \vec{A} = \frac{1}{s} \frac{\partial}{\partial s} s A_s + \frac{1}{s} \frac{\partial}{\partial \phi} A_\phi + \frac{\partial}{\partial z} A_z$

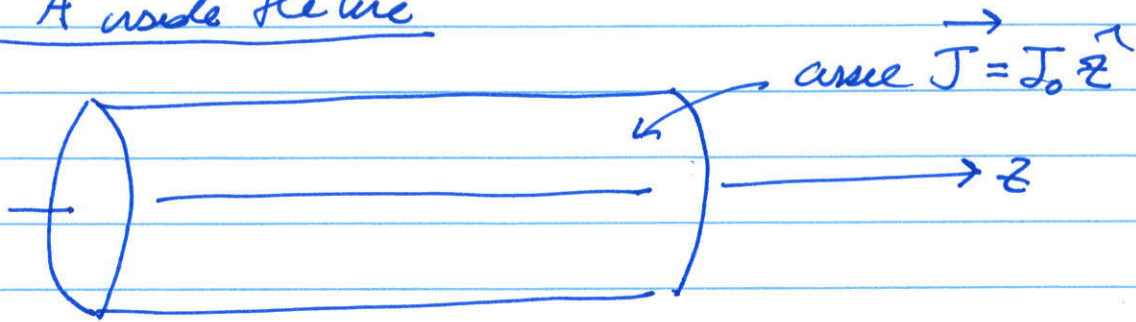
$\rightarrow 0$ , b/c  $A_s = A_\phi = 0$  and  $A_z$  does not depend on  $z$

(iv) find  $\vec{\nabla} \times \vec{A} = \hat{s} \left[ \frac{1}{s} \frac{\partial A_z}{\partial \phi} \right] + \left[ -\frac{\partial A_z}{\partial s} \right] \hat{\phi}$

$= - \left[ \frac{-\mu_0 I}{2\pi} \frac{1}{s} \right] \hat{\phi}$

$= \frac{\mu_0 I}{2\pi s} \hat{\phi}$

Sol  $\vec{A}$  inside the wire



Inside wire  $\vec{B} = \frac{\mu_0 J_0 \pi S^2}{2\pi S} \hat{\phi} = \mu_0 \frac{J_0 S}{2} \hat{\phi}$

from Ampere's law

Sol  $\vec{A}$

$$\int \vec{B} \cdot d\vec{S} = \int \vec{A} \cdot d\vec{l}$$

$$\int \mu_0 \frac{J_0}{2} S ds dz = A_z(0)l - A_z(s)l$$

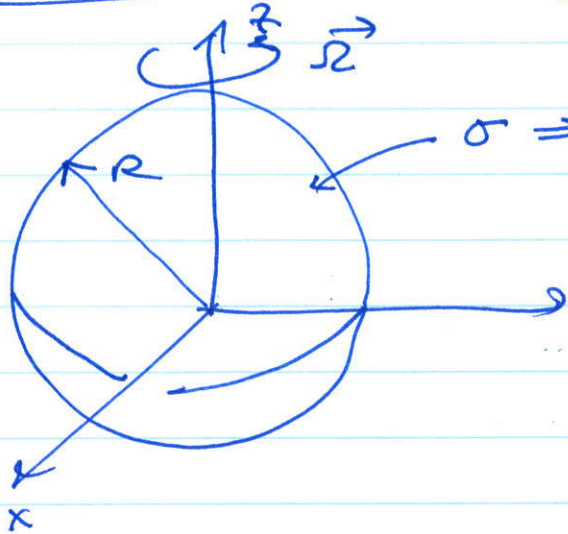
$$\mu_0 \frac{J_0 S^2}{4} \Big|_0^S l = A_z(0)l - A_z(s)l$$

$$A_z(s) = -\mu_0 \frac{J_0}{4} S^2 \frac{1}{z}$$



5.30

Find the field inside a uniformly charged sphere of charge

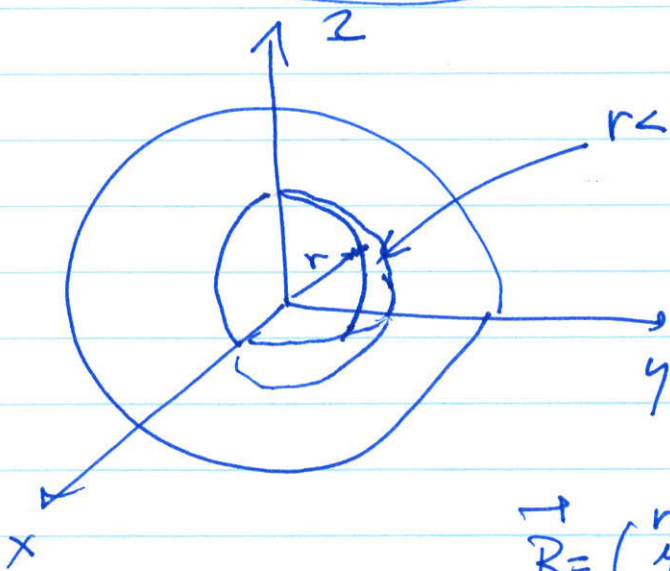


$\sigma \Rightarrow \vec{K} = \sigma \vec{\Omega} \times \vec{r} \quad | \text{ at } r=R$

④ Ex 5.11  
 $\Rightarrow \vec{B} = \frac{2}{3} \mu_0 \sigma \vec{\Omega} R$   
 inside a shell

⑥  $\vec{B} = \nabla \times \vec{A}$  outside shell  
 $= \frac{\mu_0 R^4 \Omega \sigma}{3r^3} (2\cos\theta \hat{r} + \sin\theta \hat{\theta})$

note:  $\sigma = \frac{Q}{4\pi R^2}$



$r < R$ . Need to solve 2 integrals, one for  $r' < r$  and one for  $r' > r$

~~$\vec{B} = \int \frac{\mu_0 \vec{K} \times \vec{r}'}{r^2} d\tau'$~~

$\vec{B} = \int_0^r \frac{\mu_0 r'^4 \Omega \rho dr'}{3r^3} [2\cos\theta \hat{r} + \sin\theta \hat{\theta}]$   
 $+ \int_r^R \frac{2}{3} \mu_0 \rho dr' \Omega \hat{z} r'$