

From the Weak Bruhat

Order to Crystal Posets

Patricia Hersh

North Carolina

State University

(joint work with Cristian  
Lenart, SUNY Albany)

## Perspective & Main Goal:

- Study crystal graphs regarded as posets via poset map to weak Bruhat order, namely via the (right) key map.

## Poset Structure for Many Crystals

$$u <_{\text{crystal}} v \iff u \xrightarrow{f_i} v \text{ for some } i$$

- Transfer properties of weak Bruhat order to crystals.

## Weak Bruhat order:

$$u <_{\text{weak}} s_i u \text{ if } l(s_i u) > l(u)$$

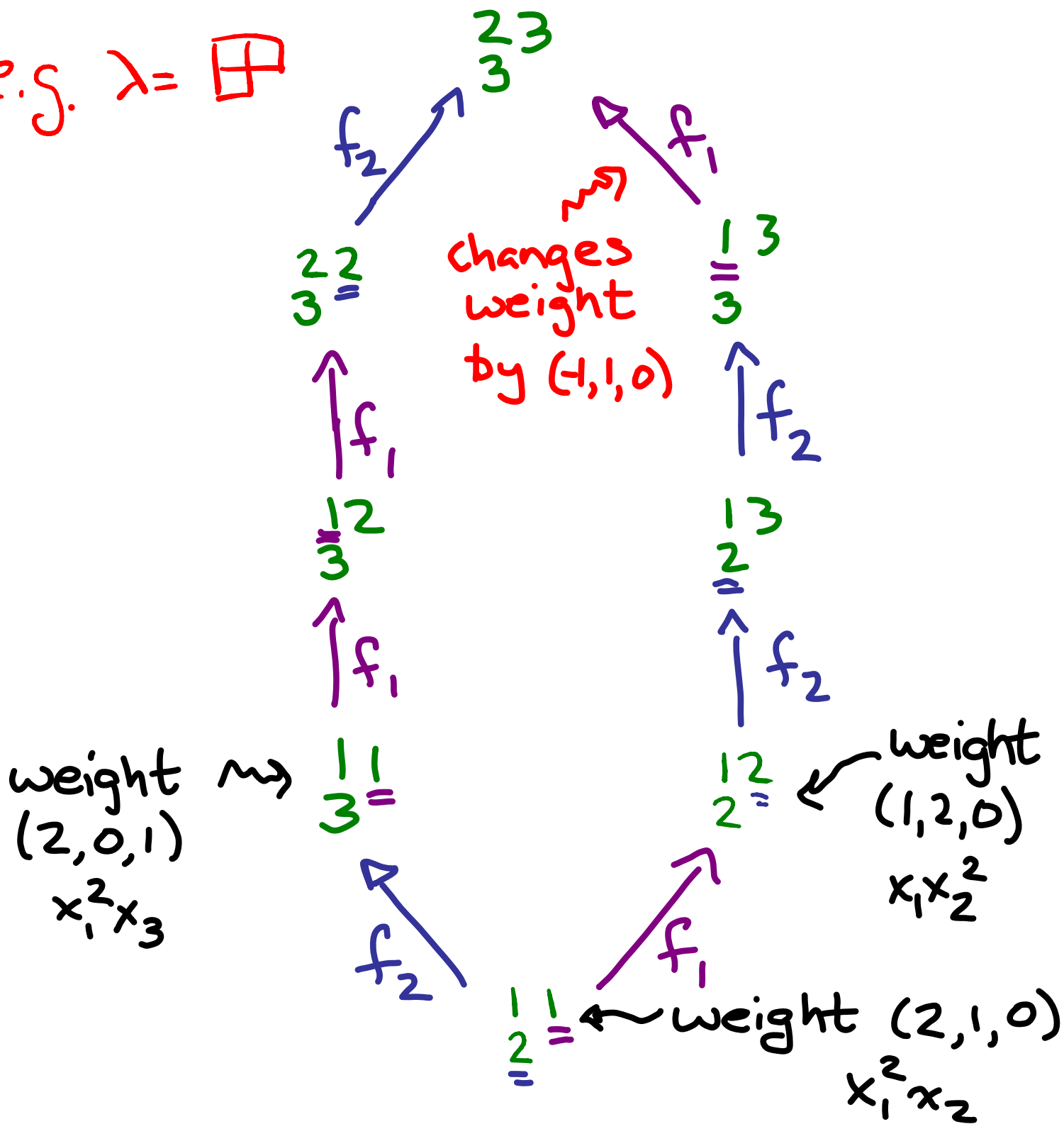
- Discover surprising new reln's amongst crystal operators

# Motivations for Crystals

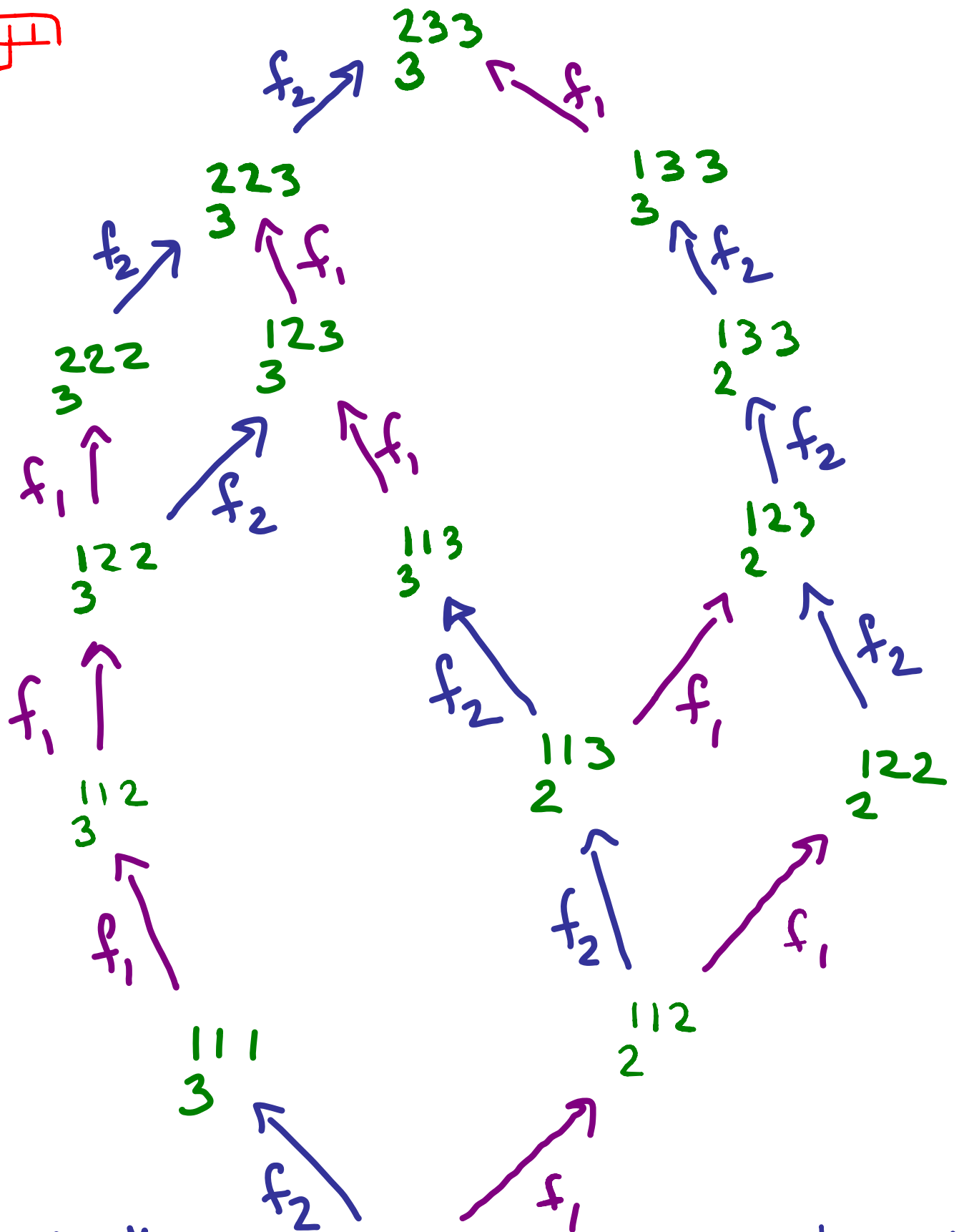
- Study representation theory of Kac-Moody algebras (e.g. affine Lie algebras)
- Take universal enveloping algebra,  $\neq$  its quantum algebra w/ parameter  $q$
- $q \rightsquigarrow 1$  yields  $U(A)$  for Kac-Moody algebra  $A$
- $q \rightsquigarrow 0$  yields algebra with same dimensions of weight spaces encoded by combinatorics of "crystal graphs" (which are often posets)
- poset elts  $\leftrightarrow$  basis vectors for weight spaces (with remarkable properties)
- cover rel's  $\leftrightarrow$  Kazhdara raising operators

# (Type A) Crystals of Highest Weight Representations & their Kashiwara Lowering Operators

e.g.  $\lambda = \begin{array}{|c|} \hline \square \\ \hline \end{array}$



$$\lambda = \begin{array}{|c|c|c|} \hline & & \\ \hline \end{array}$$



"character"  
of crystal

$$= x_1^3 x_2 + x_1^2 x_2^2 + \dots = \text{weight} \left( \begin{array}{|c|c|c|} \hline 1 & 1 & 1 \\ \hline 2 & & \\ \hline \end{array} \right) + \text{weight} \left( \begin{array}{|c|c|c|} \hline 1 & 1 & 2 \\ \hline 2 & & \\ \hline \end{array} \right) + \dots$$

highest wt vector  $(3,1,0)$  character of rep'n

# Type A crystal for highest weight rep'n of type $\lambda$

1.  $\hat{G} = \begin{matrix} 1 & 1 & 1 & - & 1 \\ 2 & 2 & - & 2 \\ 3 & 3 & - & \\ \vdots & & & \end{matrix}$  of shape  $\lambda =$  "highest weight vector"

2.  $u \xrightarrow{i} v$  for  $v$  obtained from  $u$  by incrementing to an  $i+1$  the rightmost  $i$  that is not in a "parenthesization pair" w/ an " $i+1$ "

Parenthesization Pairs: Read leftmost column bottom to top, then subsequent columns L to R, ignoring all but  $i+1$ ; pair up consec.  $i+1, i$ ; delete; repeat...

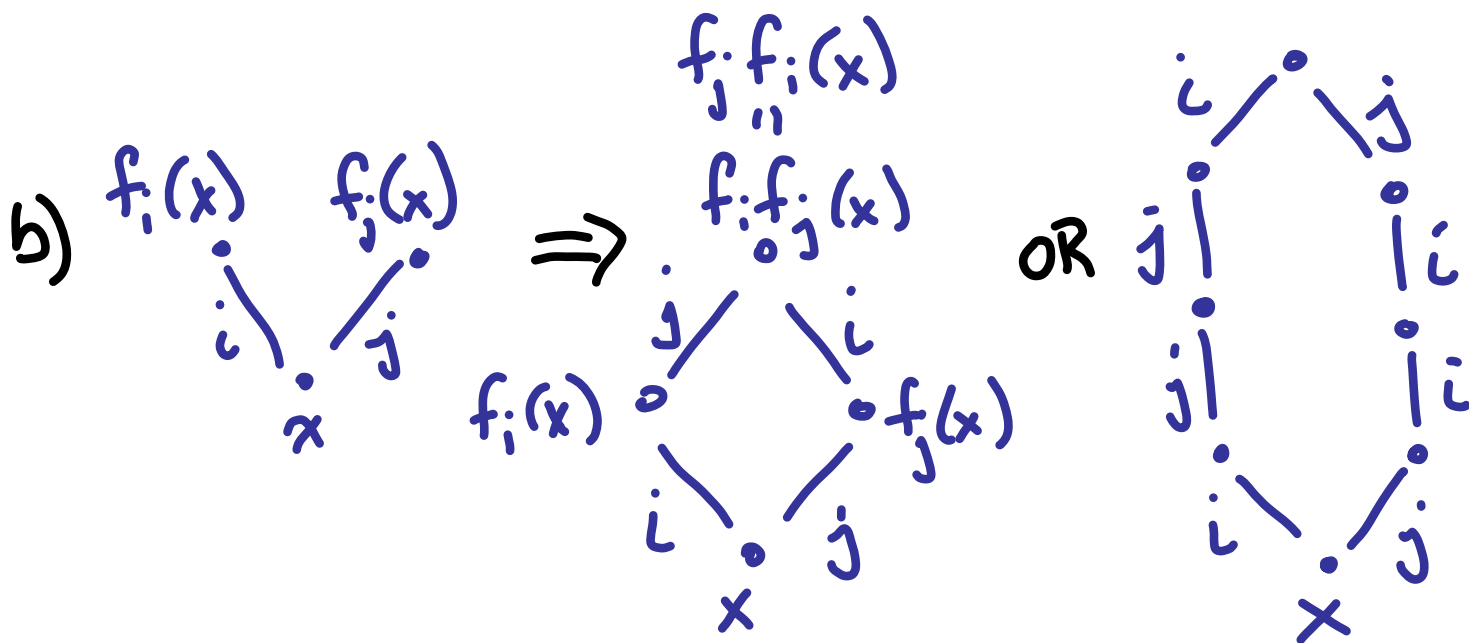
e.g.  $\begin{matrix} 1 & 1 & 1 & 1 & 4 & 4 & 4 \\ 2 & 2 & 3 & 3 \\ \boxed{3} & 4 & 4 \end{matrix}$

$\begin{matrix} 4 & 4 & 4 & 3 & 3 & 4 & 4 & 4 \\ \uparrow f_3 \\ \boxed{3} & 4 & 4 & 3 & 3 & 4 & 4 & 4 \\ \sim \rightarrow \\ i=3 \end{matrix}$

# Stanbridge Crystals: "g-crystals"

(Crystals of highest weight reps in simply laced case)

a)  $\chi_{B(\lambda)}(t) = \sum_{b \in B(\lambda)} t^{\text{wt}(b)}$  = character of irrep  $B(\lambda)$



c) likewise for  $e_i, e_j$  operators

d) axioms yield this & characterize crystals of highest weight reps in simply laced case

# Right key "k" of a KM-crystal

$$k \begin{pmatrix} 124 \\ 23 \end{pmatrix} = s_3 s_2 s_1 s_2$$



$$k \begin{pmatrix} 114 \\ 23 \end{pmatrix} = s_3 s_2 s_1$$

$$k \begin{pmatrix} 123 \\ 23 \end{pmatrix} = s_2 s_1 s_2$$



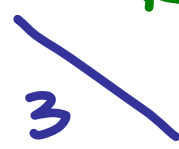
$$\begin{pmatrix} 114 \\ 22 \end{pmatrix}$$

$$k \begin{pmatrix} 113 \\ 23 \end{pmatrix} = s_2 s_1$$

$$k \begin{pmatrix} 122 \\ 23 \end{pmatrix} = s_1 s_2$$



$$k = s_3 s_2 s_1$$



$$k \begin{pmatrix} 113 \\ 22 \end{pmatrix} = s_2 s_1$$

$$k \begin{pmatrix} 112 \\ 23 \end{pmatrix} = s_1 s_2$$



$$k \begin{pmatrix} 112 \\ 22 \end{pmatrix} = s_1$$



$$k \begin{pmatrix} 111 \\ 23 \end{pmatrix} = s_2$$



$$\begin{pmatrix} 111 \\ 22 \end{pmatrix}$$

k: crystal poset  $\rightarrow$  weak Bruhat order

$$u \leq v \Rightarrow k(u) \leq k(v)$$

$$\text{Key } k(\hat{0}) = e$$



# New Algorithm to Calculate Right Key of a KM-Crystal

- (1)  $\text{key}(\hat{\sigma}) = e$
- (2) if  $\hat{\sigma} \xrightarrow{i} a$ , then  $\text{key}(a) = s_i$   
(i.e.  $\hat{\sigma} \leftarrow a$ )
- (3) if  $v$  covers 2 or more elements  
then  $\text{key}(v) = \bigvee_{\{u \mid u \rightarrow v\}} \text{key}(u)$   
(for join taken in weak order)
- (4) if  $u \xrightarrow{i} v$  and  $v$  does not cover  
any other elements, then:
  - (a)  $\text{key}(v) = \text{key}(u)$  if  $\exists u' \xrightarrow{i} u$
  - (b)  $\text{key}(v) = s_i \cdot \text{key}(u)$  otherwise

## Key Polynomials $\neq$ right / left key

(see Lascoux-Schutzenberger  $\neq$  e.g.  
Reiner-Shimozono)

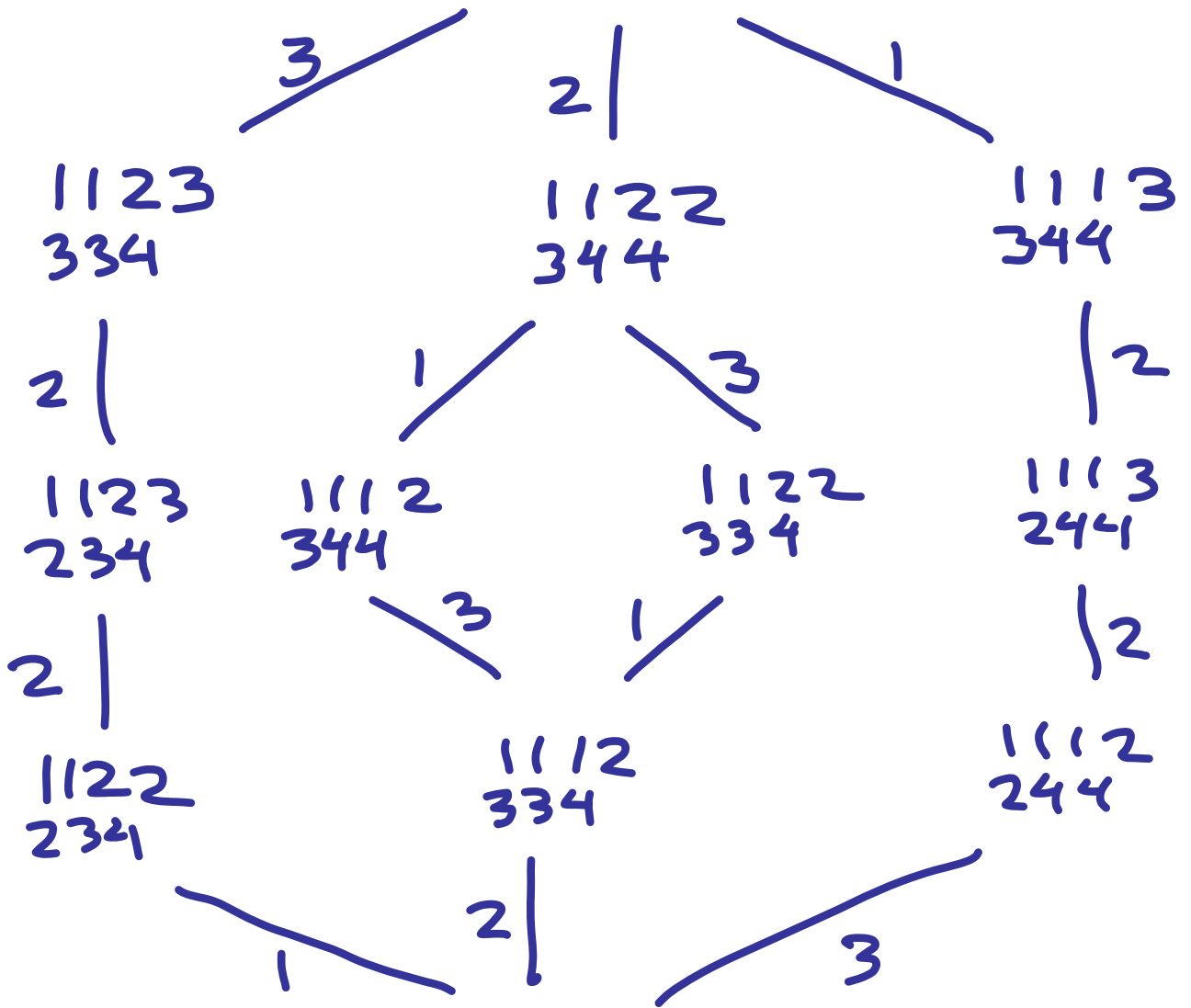
- Motivations:
- (1) Schubert poly.  $G_w$  is positive sum of "key polynomials"
  - (2) Key polynomial records character for Demazure module
  - (3) The (closely related) right  $\neq$  left key maps determine smallest Demazure modules containing a given crystal element
  - (4) These will give us poset map from  $g$ -crystal to weak Bruhat order, transferring properties

# Infinite Families of (Negative)

Examples

"Base Case":

$$v = \begin{matrix} 1123 \\ 344 \end{matrix}$$



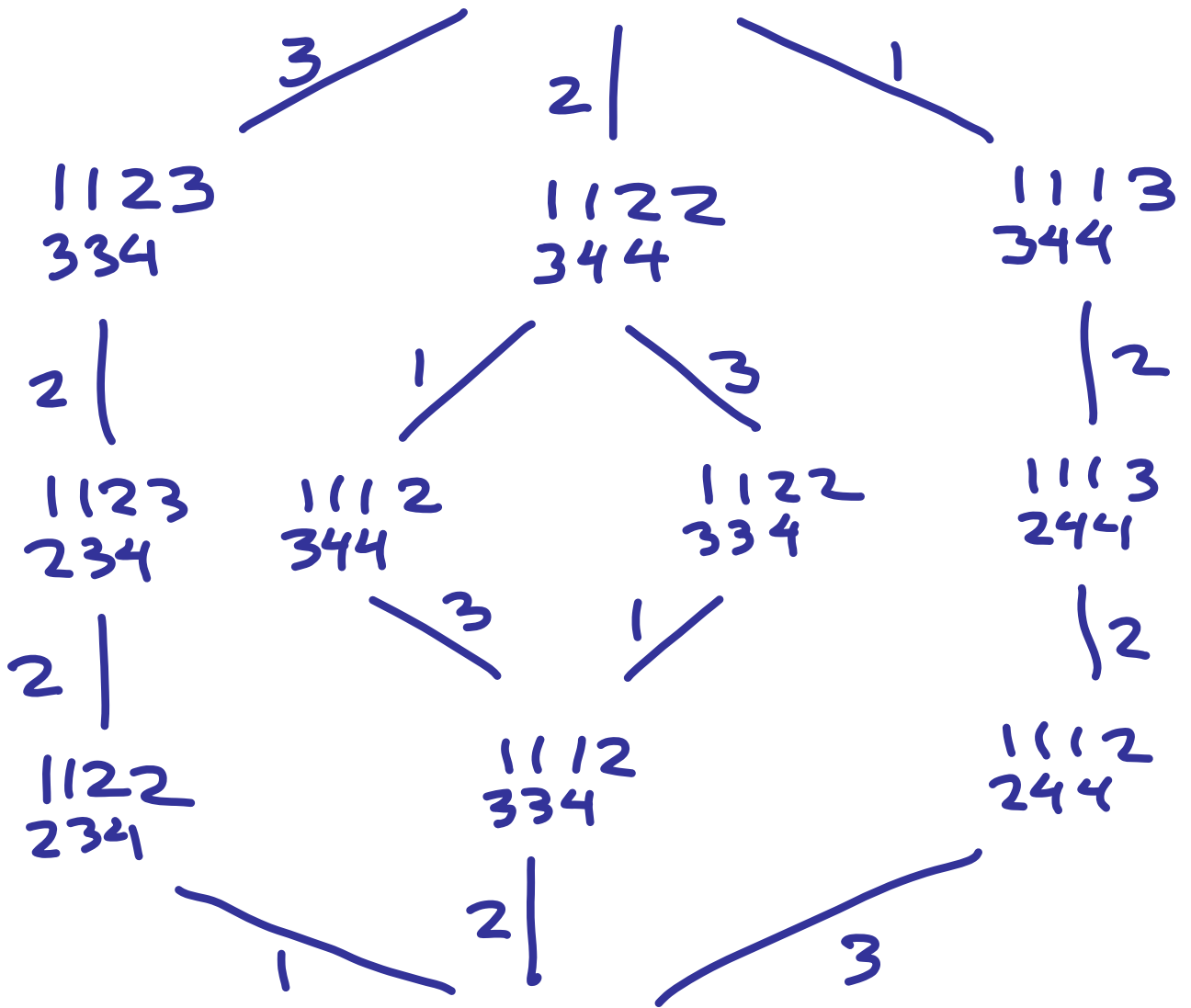
‡ not connected by "Stanbridge moves"

# Infinite Families of (Negative)

Examples

"Base Case":

$$v = \begin{matrix} 1123 \\ 344 \end{matrix}$$



$$M_p(u, v) = 2 \quad u = \begin{matrix} 1112 \\ 234 \end{matrix}$$

‡ not connected by "Steinbridge moves"

# Talk Outline (for 2nd Half)

## I. Background Review (cont.)

## II. Positive Results for Lower Intervals $[\hat{0}, u]$

- Möbius function  $\neq$  homotopy type
- Connectedness of saturated chains under "Stembridge moves"

## III. Negative Results for Arbitrary Type A Intervals $[u, v]$

- Arbitrarily large Möbius functions
- Arbitrarily high degree non-redundant "relations" amongst crystal operators
- $M(u, v) \neq q, \pm 1 \Rightarrow$  rel'n within  $[u, v]$  not generated by Stembridge local rel'n's (based on "SB-labelings")

# I. Background

Def'n: The (left) weak Bruhat order

on Coxeter system  $(W, S)$  is the partial order with cover relations

$u < \cdot v \iff v = s_i u$  for  $u, v \in W$  with

$l(v) > l(u)$ , for  $l(v) = \min \{ r \mid v = s_{i_1} \dots s_{i_r} \}$   
for  $s_{i_1}, \dots, s_{i_r} \in S$

e.g.  $W = S_n$

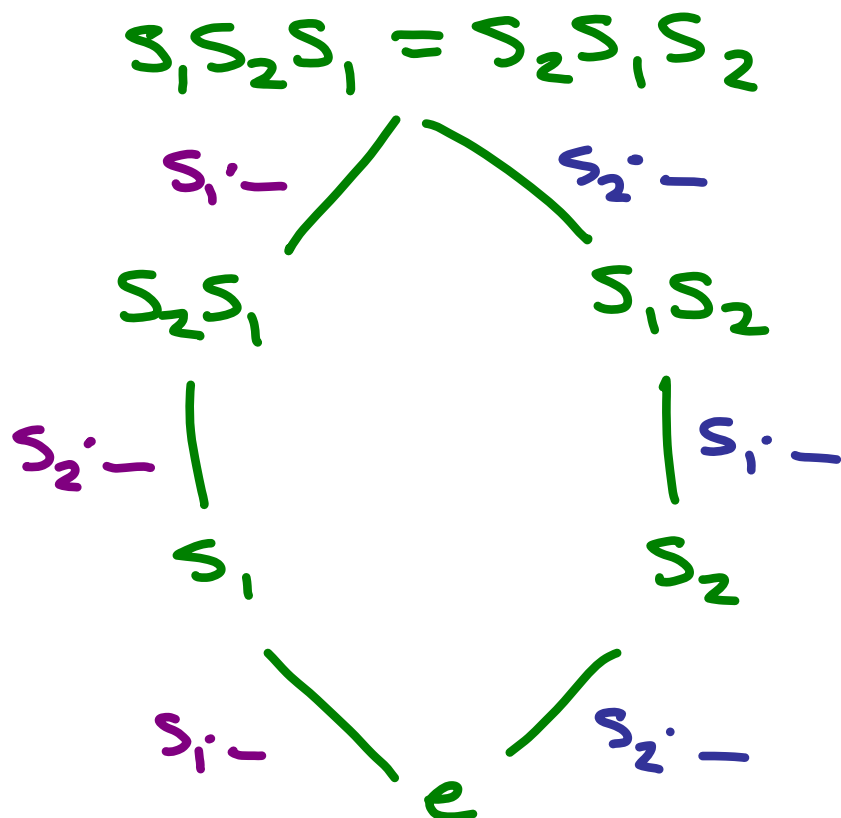
$S = \{s_1, s_2, \dots, s_{n-1}\}$  for  $s_i = (i, i+1)$

with relations:

$$s_i^2 = e \quad \& \quad s_i s_{i+1} s_i = s_{i+1} s_i s_{i+1} \quad \& \quad s_i s_j = s_j s_i \\ \text{(for } |j-i| > 1 \text{)}$$

"braid rel's"

e.g. Left weak order for  $S_3$



Key Fact: Saturated chains from  $e$  to  $w$  naturally labeled w/ the "reduced expressions"  $s_{i_1} \dots s_{i_{\ell(w)}}$  for  $w$ .  
Likewise, saturated chains from  $u$  to  $v \iff$  reduced expressions for  $v u^{-1}$

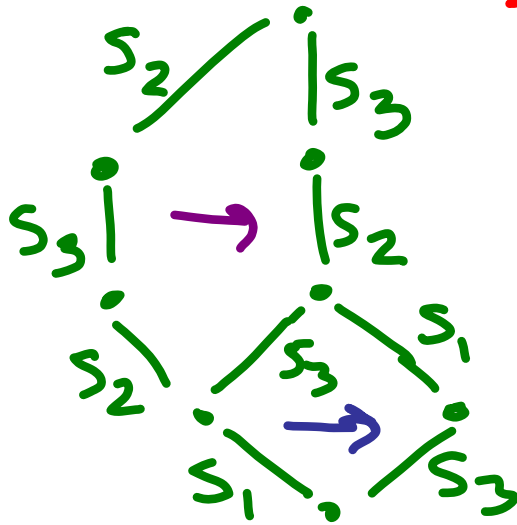
# Connectedness under Braid Moves

Thm 3.3.1 (in Björner-Brenti) Let  $(W, S)$  be a Coxeter group  $w/w \in W$ . Then every two reduced expressions for  $w$  are connected via braid moves.

e.g.  $s_2 s_3 s_2 s_1 \rightarrow s_3 s_2 s_3 s_1$   
 $\rightarrow s_3 s_2 s_1 s_3$

$w = s_2 s_3 s_2 s_1 = s_3 s_2 s_3 s_1 = s_3 s_2 s_1 s_3$

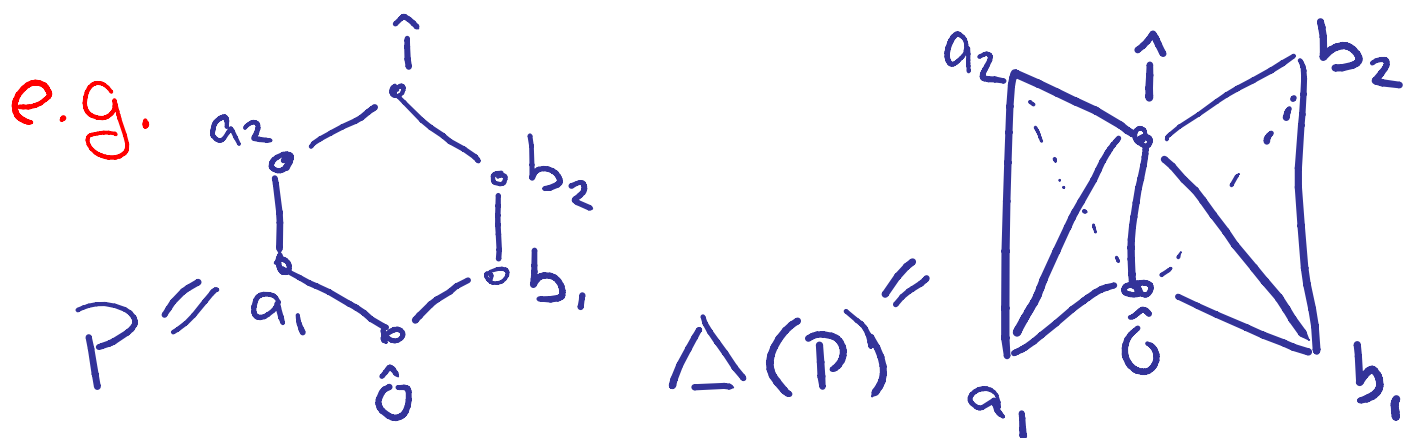
Right weak order:



Note: Proof via lattice property for  $[u, v]$



Def'n: The **order complex** (or **nerve**) of a poset  $P$  is the simplicial complex  $\Delta(P)$  whose  $i$ -dimensional faces are the  $(i+1)$ -chains  $v_0 < \dots < v_i$  in  $P$



Recall:  $M_P(u, v) = \tilde{\chi}(\Delta(\underbrace{u, v}_{\{z \in P \mid u < z < v\}}))$

( $M_P(u, v) = 0, \pm 1$  suggests ball or sphere)

## II. Positive Results for Lower Intervals $[\hat{0}, u]$

Recall:  $M_p(u, u) = 1$

$$M_p(u, v) = -\sum_{u \leq z < v} M_p(u, z)$$

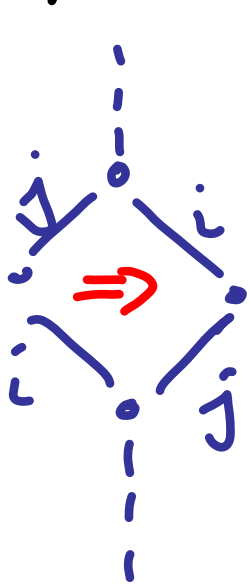
Thm 1 (H.-Lenart): Given  $u$  in a symmetrizable Kac-Moody type crystal "KM-crystal", then  $M(\hat{0}, u) = 0, \pm 1$ . More specifically,  $M(\hat{0}, u) = 0$  unless  $\text{key}(u) = \omega_0(J)$  for some parabolic subgroup  $\omega_J$  with  $u$  the unique smallest element in  $\text{key}^{-1}(\omega_0(J))$ , in which case  $M(\hat{0}, u) = (-1)^{|J|}$ .

Thm 2 (H.-Lenart): Given a symmet. KM-crystal  $\dagger$  given any parabolic  $w_J$ , then  $\text{key}^{-1}(w_0(w_J))$  has a unique minimal element and a unique maximal element.  
 (Proof via alcove path model)

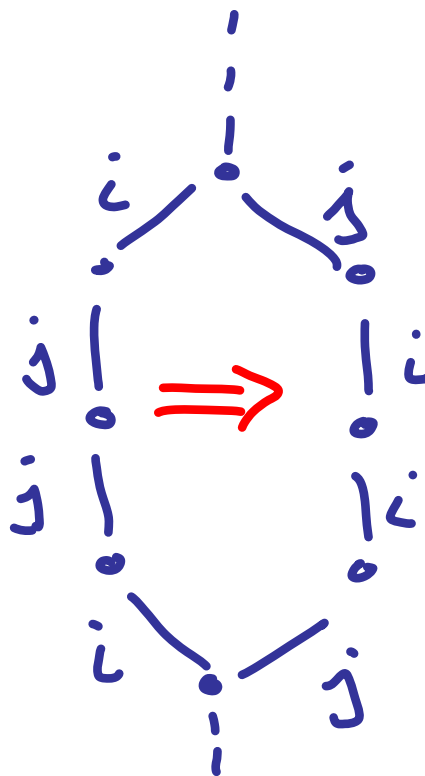
Thm 3 (H.-Lenart): Each lower interval  $(\hat{0}, u)$  in a symmet. KM-crystal has  $\Delta(\hat{0}, u) \simeq$  ball or sphere, getting  $S^{|\mathbb{J}|-2}$  for  $u = \min(k^{-1}(w_0(\mathbb{J})))$ . Likewise for upper intervals in finite KM-crystals.  
 (Proof via Quillen fibre lemma)

Thm 4 (H.-Lenart): Given any lower interval  $(\hat{0}, u)$  in a  $\gamma$ -crystal, then set of saturated chains from  $\hat{0}$  to  $u$  is connected by "**Stanbridge moves**", namely moves of the

form



and

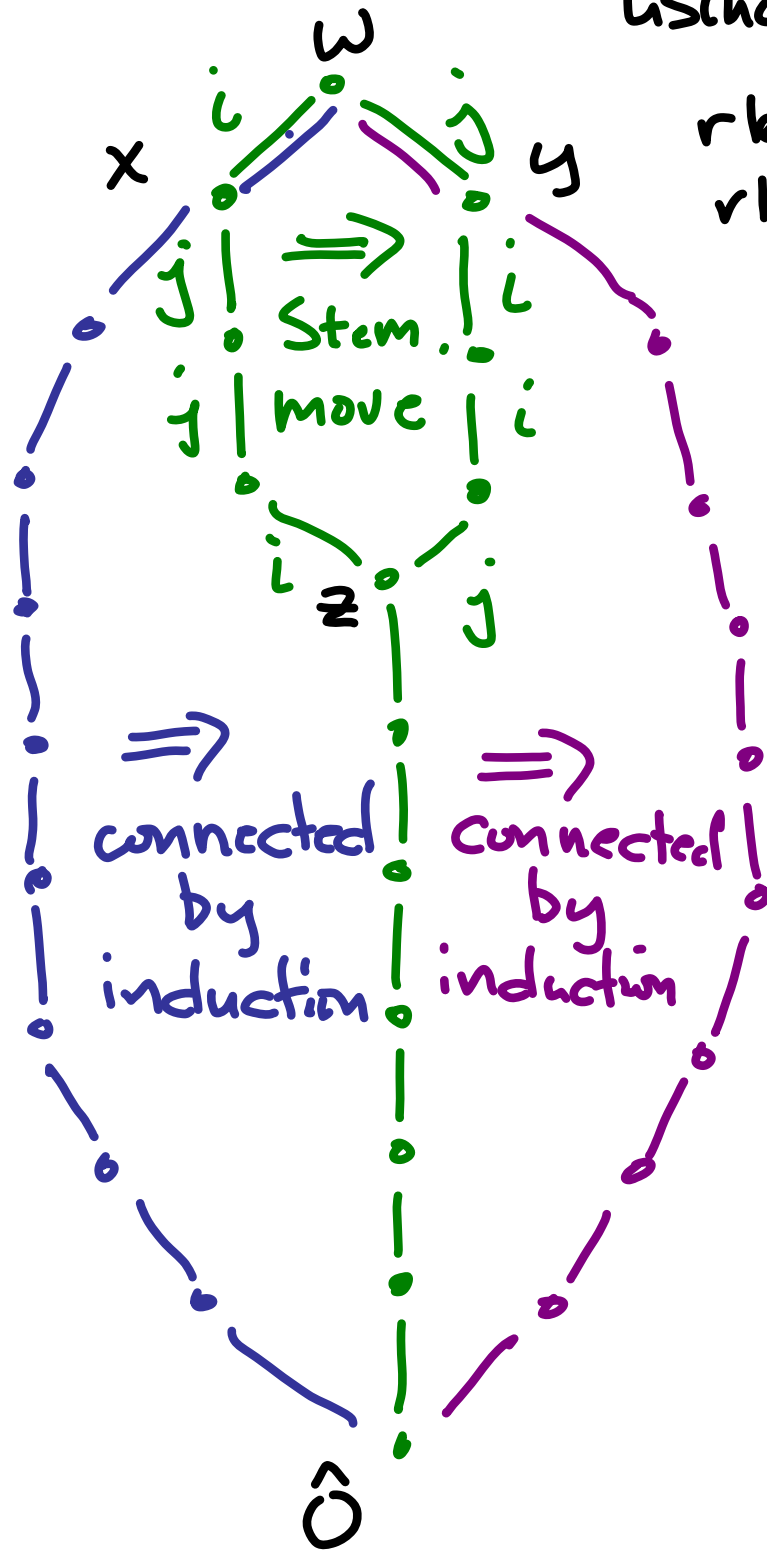


Note: Likewise in doubly-faced case via "**Stanberg moves**".

Proof Idea: Induction on rank

using  $z \geq \hat{0}$  †

$rk(x) < rk(w)$   
 $rk(y) < rk(w)$



### III. Negative Results for Arbitrary (not necessarily lower)

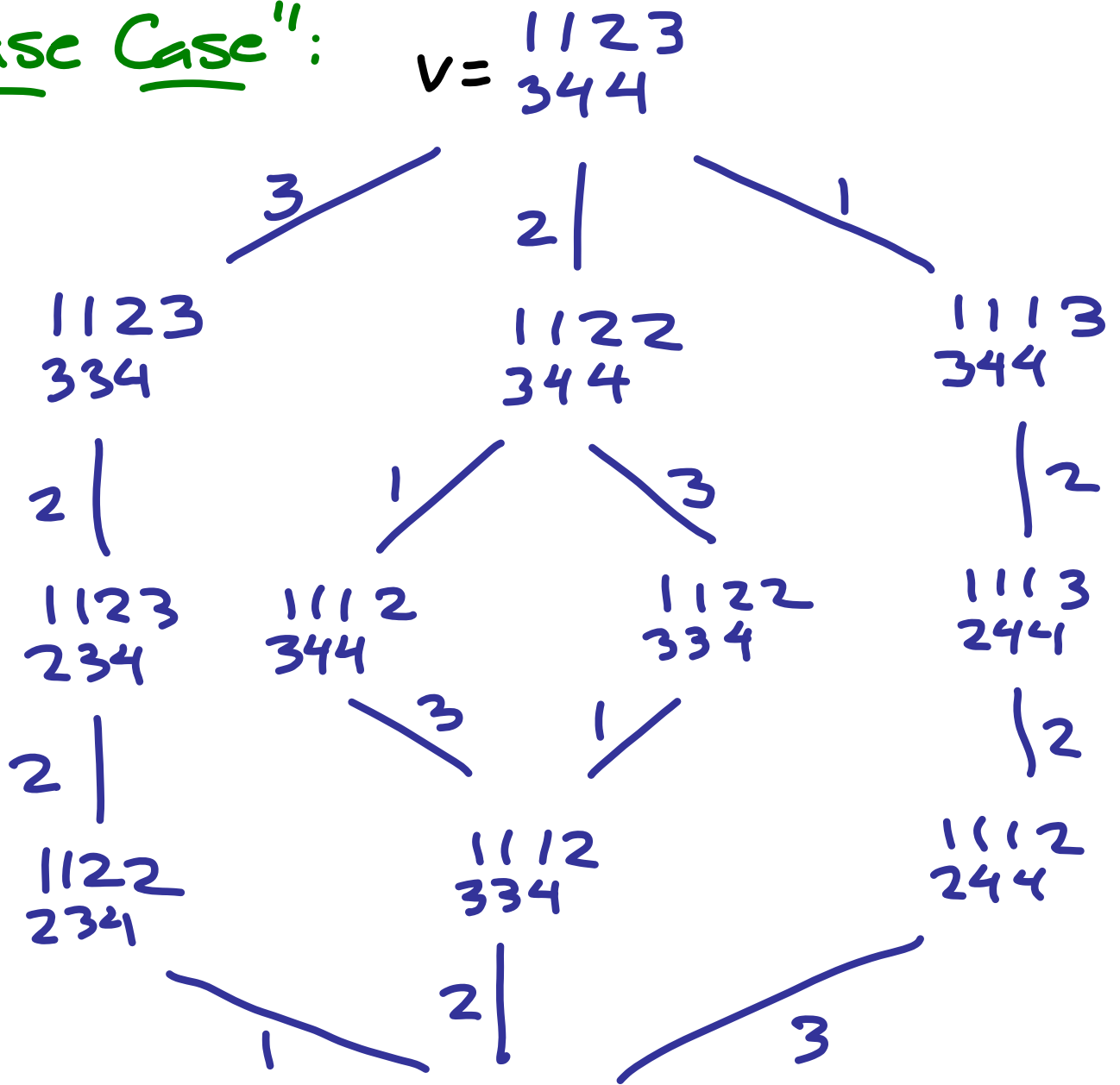
#### Crystal Poset Intervals (type A)

Thm 5 (H. - Lenart): There exist elements  $u, v$  in type A  $g$ -crystals with  $M(u, v) = 2j$  for every positive integer  $j$ .

Thm 6 (H. Lenart): There exist type A intervals  $[u, v]$  with  $\text{rk}(v) - \text{rk}(u)$  arbitrarily large s.t.  $(u, v)$  is disconnected

# Infinite Family of Examples

"Base Case":



$M_p(u, v) = 2$       $u = \begin{matrix} 1112 \\ 234 \end{matrix}$

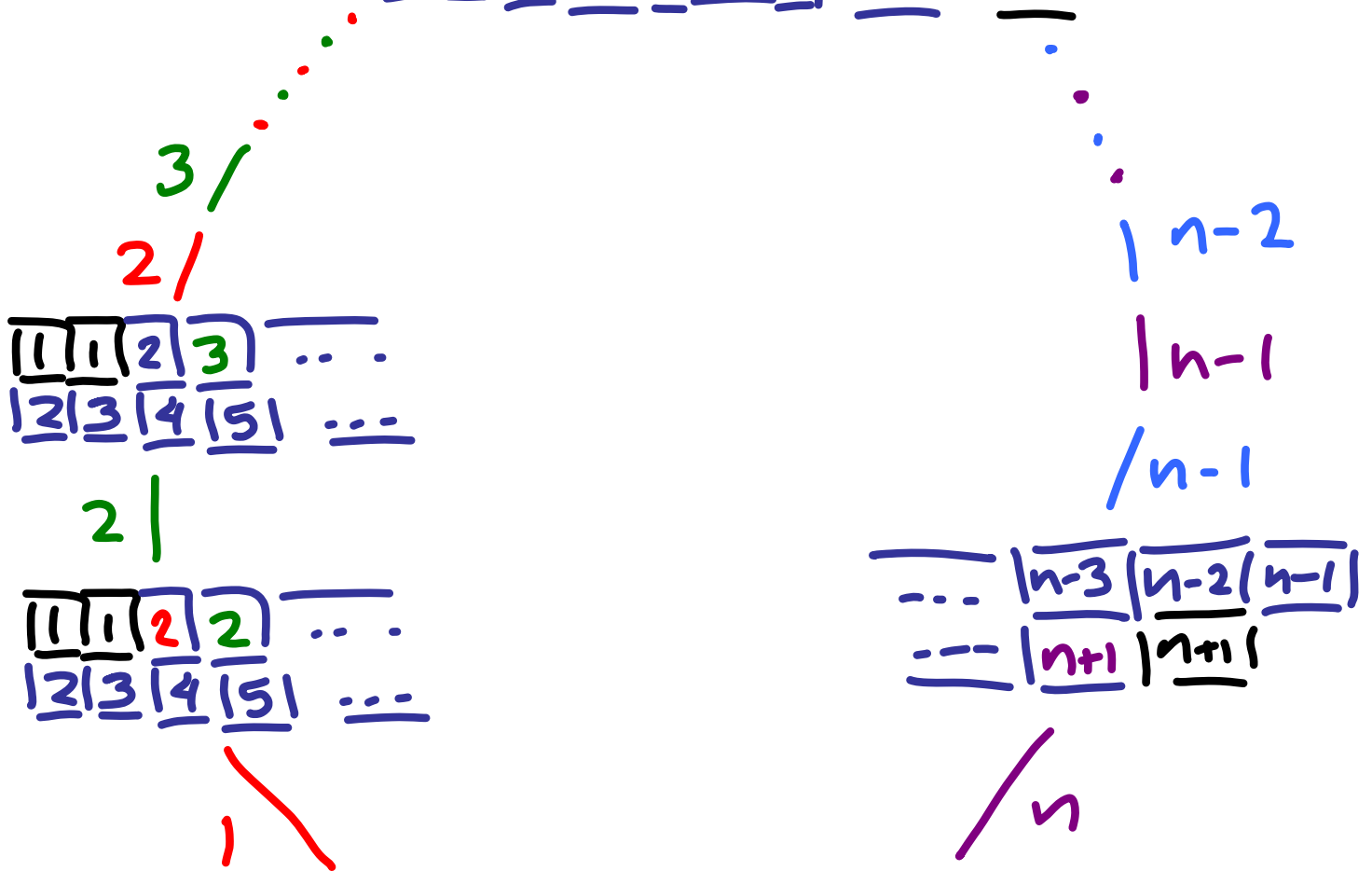
‡ not connected by "Steinbridge moves"

# Arbitrarily High Rank

## Disconnected Open Intervals

$$v = \begin{array}{|c|c|c|c|} \hline 1 & 1 & 2 & 3 \\ \hline \end{array} \dots \begin{array}{|c|c|c|} \hline n-2 & n-1 & n \\ \hline \end{array}$$

$$\begin{array}{|c|c|c|c|} \hline 3 & 4 & 5 & 6 \\ \hline \end{array} \dots \begin{array}{|c|c|} \hline n+1 & n+1 \\ \hline \end{array}$$



$$u = \begin{array}{|c|c|c|c|} \hline 1 & 1 & 1 & 2 \\ \hline \end{array} \dots \begin{array}{|c|c|c|} \hline n-3 & n-2 & n-1 \\ \hline \end{array}$$

$$\begin{array}{|c|c|c|c|} \hline 2 & 3 & 4 & 5 \\ \hline \end{array} \dots \begin{array}{|c|c|} \hline n & n+1 \\ \hline \end{array}$$

label sequences:  $1, 2, 2, 3, 3, 4, 4, \dots, n-1, n-1, n$   
 $\neq n, n, n, \dots, 2, 2, 1$  in distinct components



Consequence: Arbitrarily high degree relations  $e_{i_1} \dots e_{i_d}(u) = e_{j_1} \dots e_{j_d}(u)$  amongst crystal operators applied to  $u$  not implied by any lower degree relations.

Systematic Method to find Such Unexpected Relations?

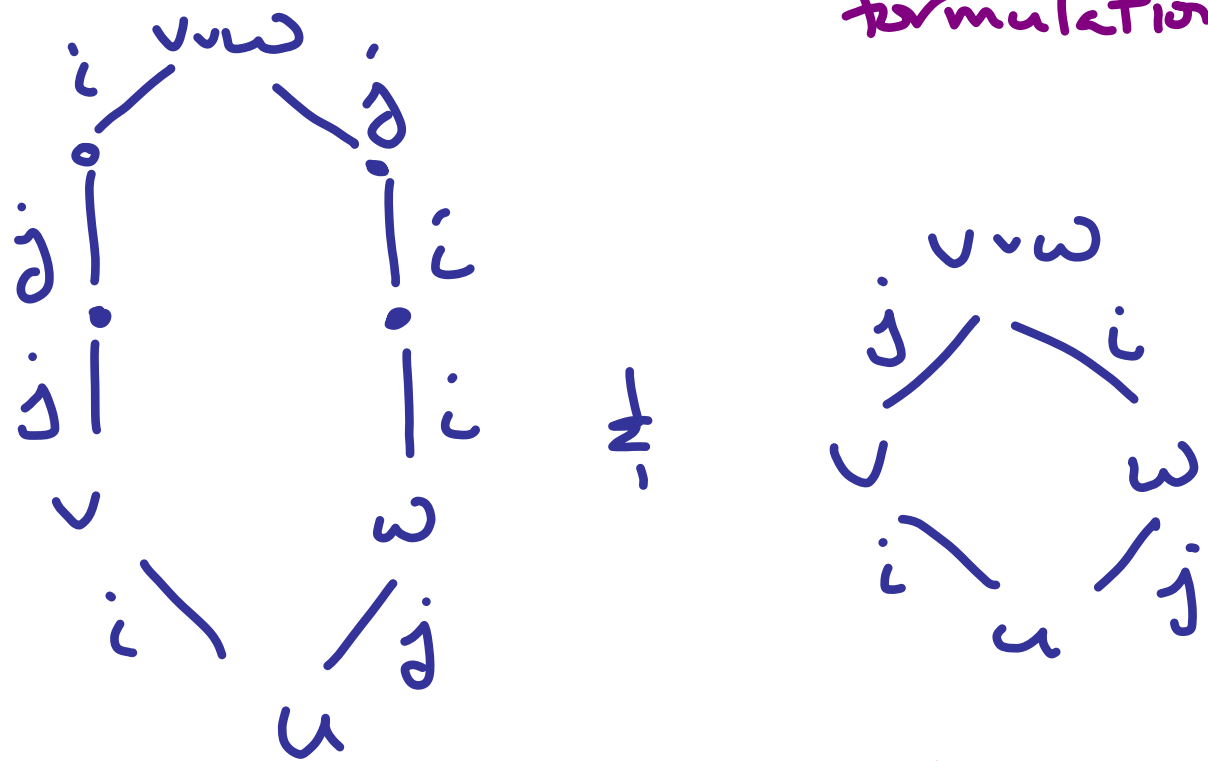
- Möbius functions, due to theory of "SB-labelings" of H.-Meszáros

Def'n (H.-Meszáros): A finite lattice  $L$  has SB-labeling if it has edge labeling  $\lambda$  s.t. for  $u < v \neq u < w$

(1)  $\lambda(u,v) \neq \lambda(u,w)$  for  $v \neq w$

(2) Each saturated chain from  $u$  to  $v \vee w$  uses both these labels  $\neq$  no others. (We call this "index 2 formulation")

e.g.



meet the requirements!

Thm (H.-Meszáros):  $\lambda$  satisfies index 2 formulation for SB-labeling  $\Leftrightarrow \lambda$  satisfies global formulation for SB-labeling.

Thm (H.-Meszáros): If a finite lattice  $L$  has SB-labeling, then  $M_L(x, y) = 0, \pm 1$  for all  $x, y \in L$ .

Observation (H.-Lenart): If a crystal w/ unique minimal element were a lattice s.t. Stembridge local structure gives least upper bounds for  $v \neq w$  s.t.  $u \prec v$  &  $u \prec w$ , then edge coloring would be SB-labeling.

Lemma (H.-Lenart): If crystal has  $u \prec f_i(u)$  &  $u \prec f_j(u)$  s.t.  $x = f_i f_j(u) = f_j f_i(u)$  (resp.  $x = f_i f_j^2 f_i(u) = f_j f_i^2 f_j(u)$ ), then there does not exist  $x'$  s.t.  $u \prec x' \prec x$ .

Thm (H.-Lenart): If crystal of rep'n has  $M(x, y) \neq 0, \pm 1$ , then have rel'n of crystal operators within  $[x, y]$  that is not implied by Stembridge local rel'ns.

Appendix: a few slides  
with extra details...

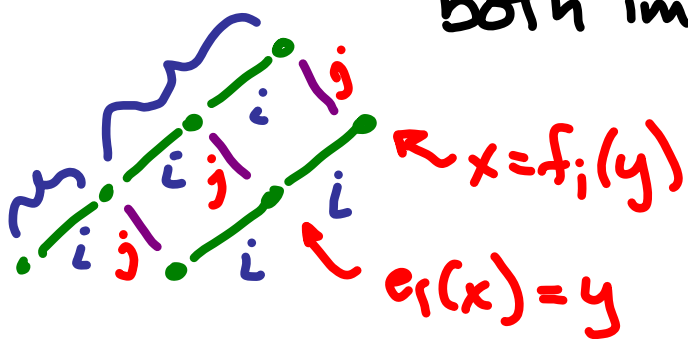
# Crystals

A **crystal**  $B$  of type  $\phi$  is a nonempty set  $B$  with raising & lowering operators  $e_i, f_i$  & maps

$$\varepsilon_i, \varphi_i : B \rightarrow \mathbb{Z} \cup \{-\infty\}$$

$\text{wt} : B \rightarrow \Lambda = \text{weight lattice of type } \phi$   
s.t.

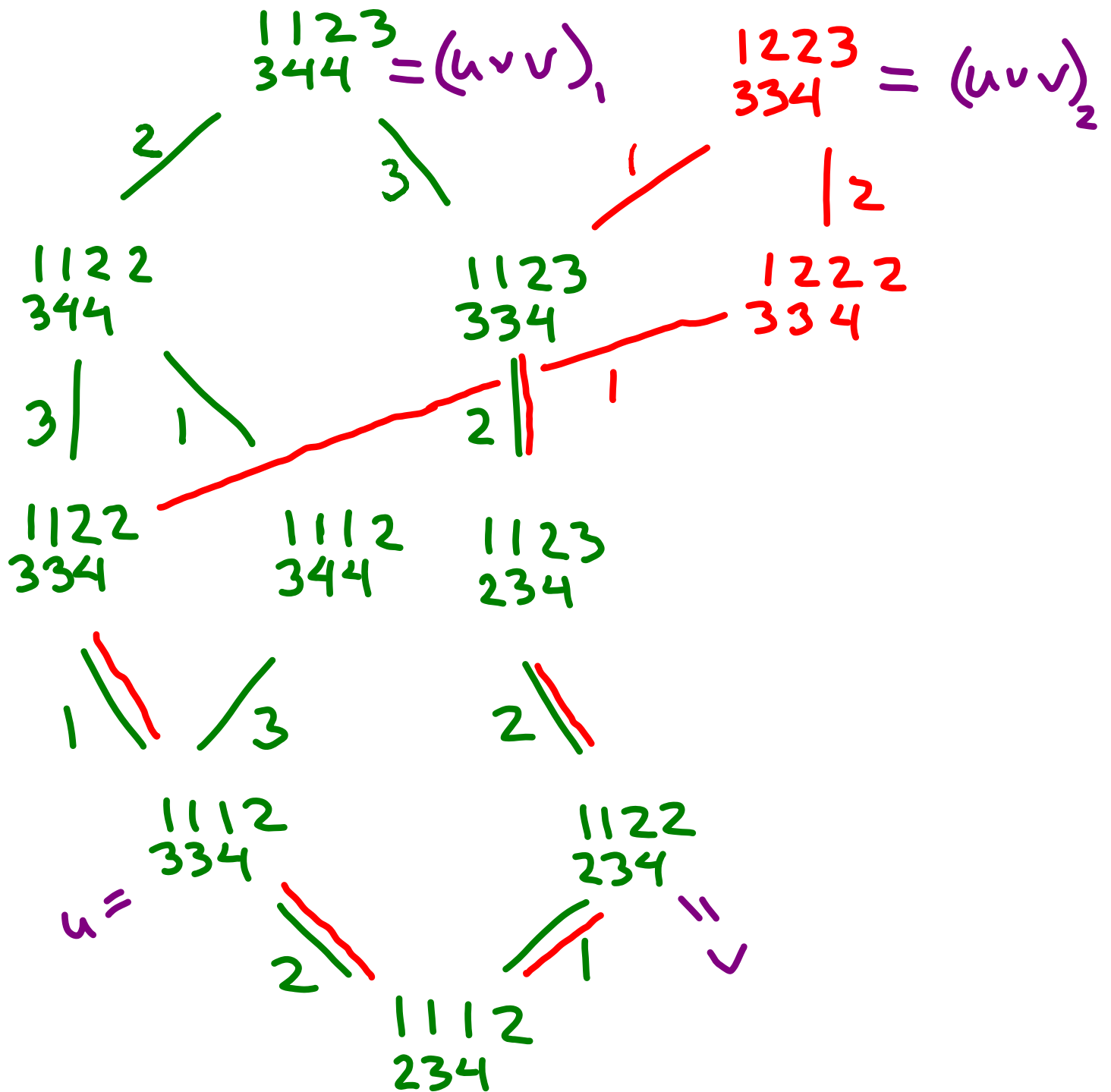
(A1)  $x, y \in B$ , then  $e_i(x) = y \Leftrightarrow x = f_i(y)$   
both implying  $\text{wt}(y) = \text{wt}(x) + \alpha_i$



$$\begin{aligned} \varepsilon_i(y) &= \varepsilon_i(x) - 1 \\ \varphi_i(y) &= \varphi_i(x) + 1 \end{aligned}$$

(A2)  $\varphi_i(x) - \varepsilon_i(x) = \langle \text{wt}(x), \alpha_i^\vee \rangle$

# Non-Lattice Example:



# Examples with $M(u,v) = 2^j$

$j=1: u = \begin{matrix} 1112 \\ 234 \end{matrix}$

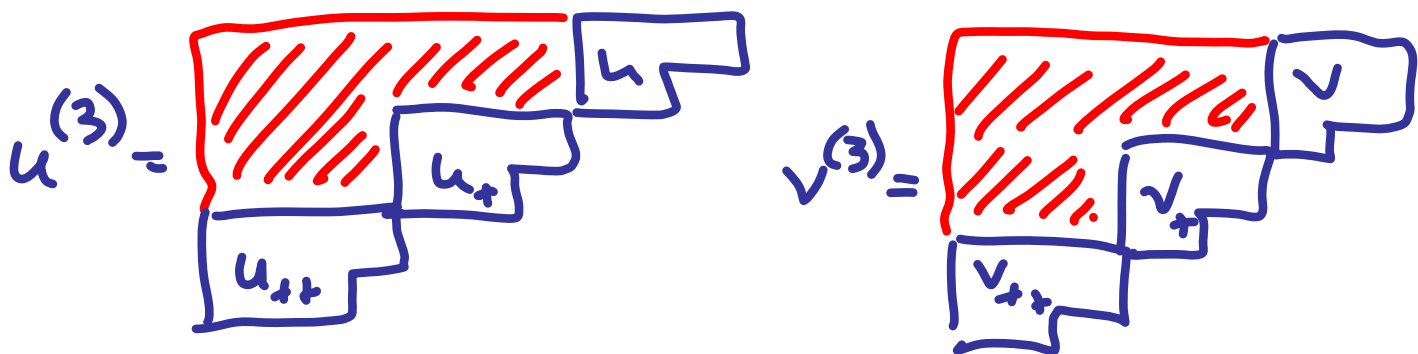
$v = \begin{matrix} 1123 \\ 344 \end{matrix}$

$j=2: u^{(2)} = \begin{matrix} 1111 & \boxed{1112} \\ 2222 & \boxed{234} \\ \boxed{6667} & \\ \boxed{789} & \end{matrix}$        $v^{(2)} = \begin{matrix} 1111 & \boxed{1123} \\ 2222 & \boxed{344} \\ \boxed{6678} & \\ \boxed{899} & \end{matrix}$

$u_+ := u+S = \begin{matrix} \boxed{6667} \\ \boxed{789} \end{matrix}$        $v_+ := v+S = \begin{matrix} \boxed{6678} \\ \boxed{899} \end{matrix}$

$[u^{(2)}, v^{(2)}] \cong [u, v] \times [u, v]$

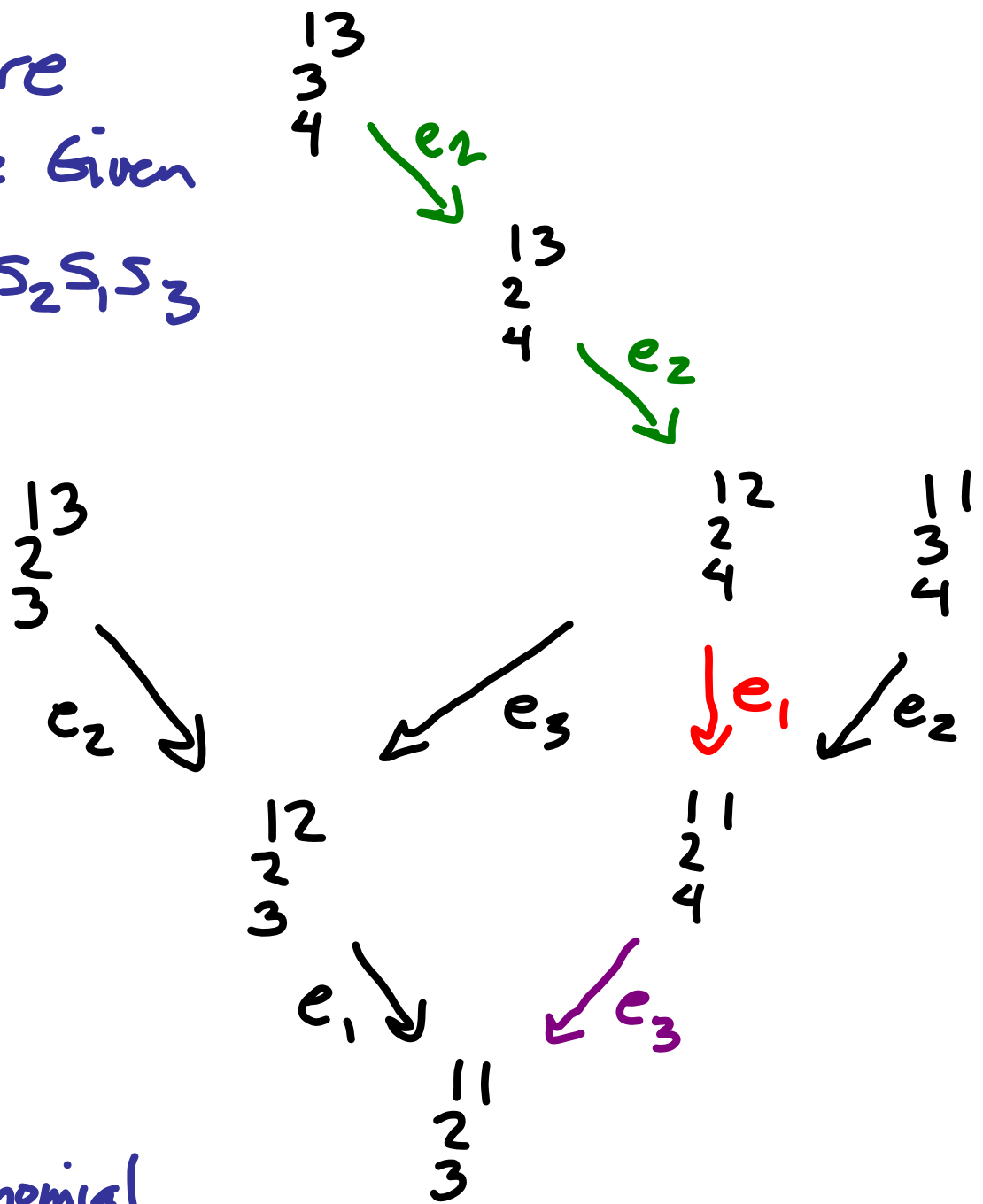
so  $M(u^{(2)}, v^{(2)}) = 2^2$



$[u^{(k)}, v^{(k)}] \cong \underbrace{[u, v] \times \dots \times [u, v]}_{k\text{-fold}} \quad M = 2^k$



Demazure  
Module Given  
by  $w = s_2 s_1 s_3$



Key Polynomial

$$K_{(1,0,2,1)} = \sum_{T' \leq T} x^{T'}$$

componentwise

$\updownarrow$   
 $K(T') \leq \text{Bruhat } K(T) \text{ w/ no higher } e_i \text{ exponents}$

# Crucial Properties of Key

Thm (Littlemann): Given any symmetrizable Kac-Moody algebra  $A$ , the key of any crystal of type  $A$  satisfies:

$$K(f_p(F)) = \begin{cases} K(F) & \text{if } e_p(F) \neq 0 \\ s_p K(F) \text{ or } K(F) & \text{if } e_p(F) = 0 \end{cases}$$

Also, if  $e_p(F) = 0$  then  $s_p K(F) > K(F)$

Corollary: If  $K(F) = s_{i_1} \dots s_{i_r}$  then there exists saturated chain from  $F$  to  $\hat{0}$  given by applying  $e_{i_r}^{d_r} \dots e_{i_1}^{d_1}$  to  $F$  for some  $d_1, \dots, d_r > 0$ .

# Relation to Reiner-Shimozono Viewpoint on Key Polynomials

- $\partial_i = \frac{1-s_i}{x_i - x_{i+1}} \quad \neq \quad \pi_i = \partial_i x_i$
- $K_\alpha = \pi_{i_1} \dots \pi_{i_r} x^{\lambda(\alpha)}$  for  $\alpha$  composition of  $n \neq s_{i_1} \dots s_{i_r}$  sorting  $\alpha$  to  $\lambda(\alpha)$

e.g.  $K_{(1,0,2,1)} = \pi_2 \pi_1 \pi_3 x^{(2,1,1,0)}$

$= \pi_2 \pi_1 (x_1^2 x_2 (x_3 + x_4))$

$= \pi_2 (x_1 x_2 x_3 (x_1 + x_2) + x_1 x_2 x_4 (x_1 + x_2))$

$= x_1^2 x_2 x_3 + x_1 x_2 x_3 (x_2 + x_3) + x_1^2 x_4 (x_2 + x_3) + x_1 x_4 (x_2^2 + x_2 x_3 + x_3^2)$