

# SMEFT vs HEFT

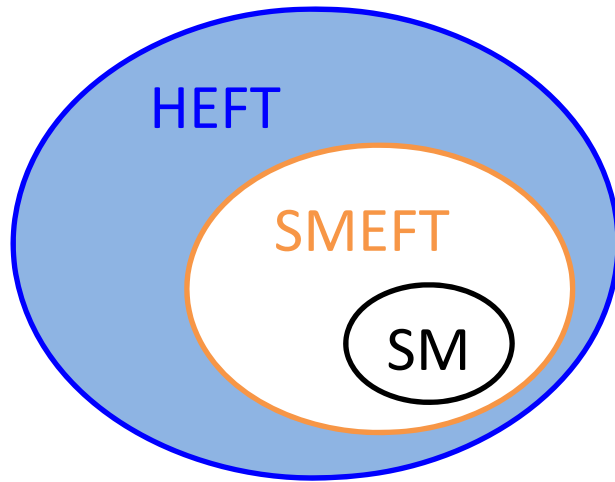
Pacific NW Theory Workshop, Sep 21-25, 2020

Xiaochuan Lu

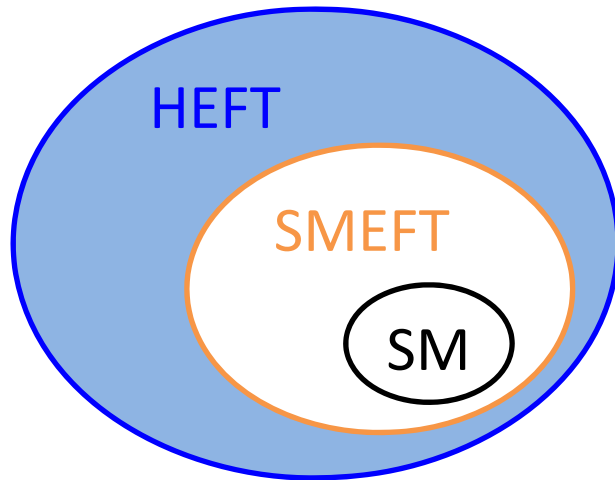
University of Oregon

arXiv: 2008.08597,

with Timothy Cohen, Nathaniel Craig, and Dave Sutherland

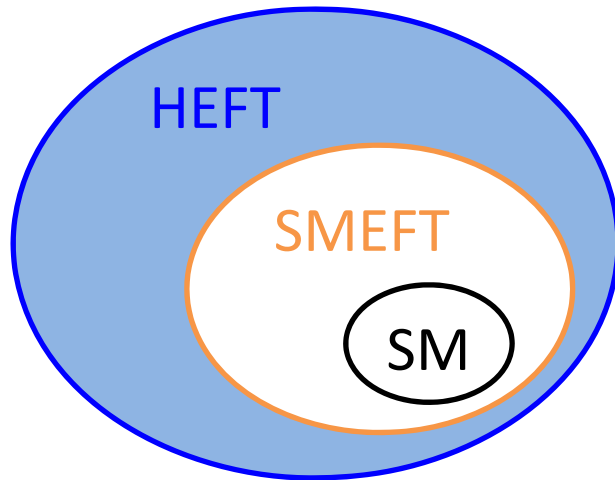


$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \frac{C_H}{\Lambda^2} |H|^6 + \frac{C_{H\Box}}{\Lambda^2} |H|^2 \partial^2 |H|^2 + \frac{C_R}{\Lambda^2} |H|^2 |DH|^2 + \dots$$



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$$H \rightarrow \left\{ h, U \equiv \exp(i\pi^a t^a / v) \right\}$$

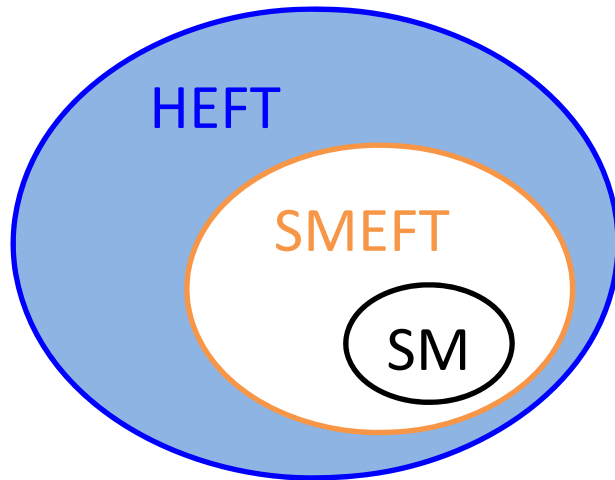


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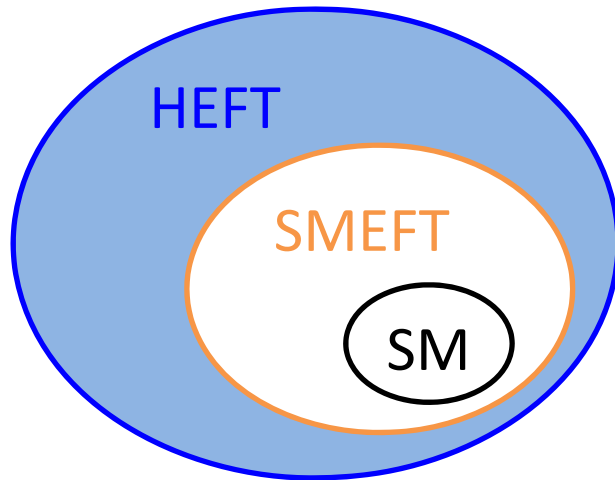
$$\begin{cases} \Sigma^\dagger \Sigma = |H|^2 \mathbf{1}_{2 \times 2} \\ \det(\Sigma) = |H|^2 = \frac{1}{2}(v+h)^2 \end{cases}$$



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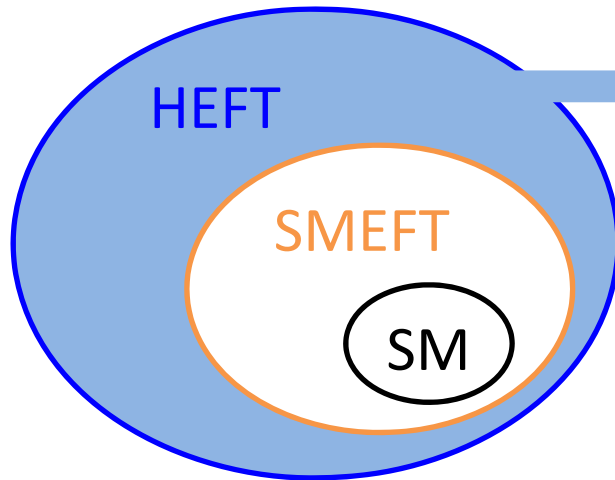


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## Outline

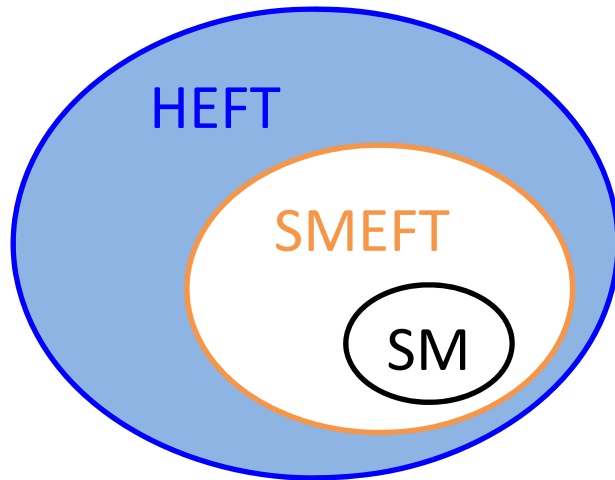
- What is HEFT?
- What UV physics generate HEFT?

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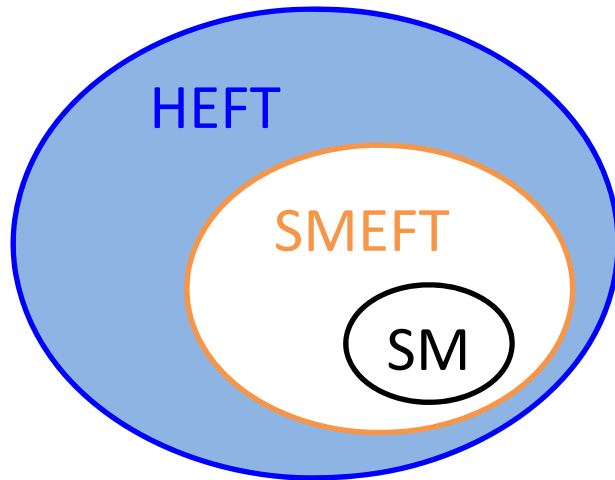


**Markus Luty**

**That is crap!**

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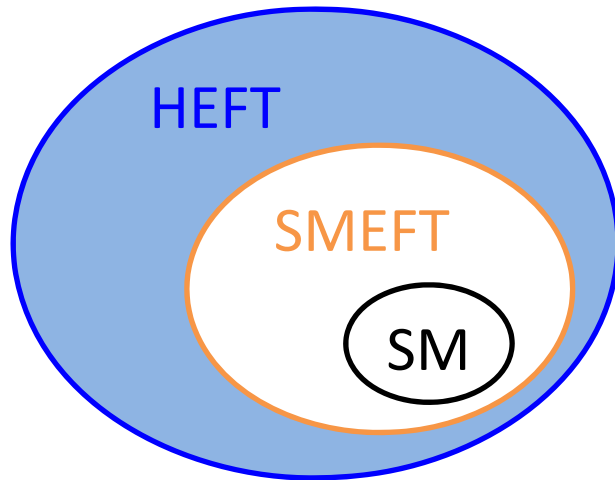


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$$\begin{cases} U = \frac{\sqrt{2}}{v+h} \Sigma \\ h = \sqrt{2|H|^2} \end{cases}$$

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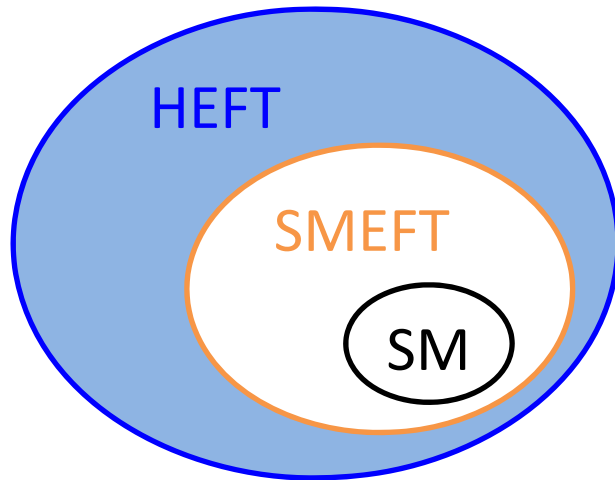
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$$\begin{cases} \text{SMEFT: } (x, y) \\ \text{HEFT: } (r, \theta) \end{cases} \Rightarrow r = \sqrt{x^2 + y^2}$$

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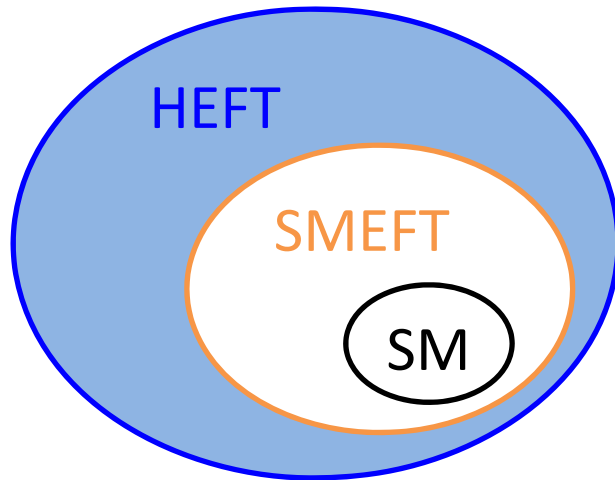
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$$\begin{aligned} \sqrt{x} &= \sqrt{x_0 + (x - x_0)} \\ &= \sqrt{x_0} + \frac{1}{2\sqrt{x_0}}(x - x_0) + \dots \end{aligned}$$

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## Adam Falkowski and Riccardo Rattazzi: (arXiv: 1902.05936)

UV cut-off. Our distinction between analytic and non-analytic lagrangians coincides with the distinction, in use in the Higgs EFT community, between linear (so-called SMEFT) and non-linear (so-called HEFT) effective theory, or equivalently between  $h$  being or not being part of a  $SU(2)_W$  doublet. We however believe our classification is more adequate and enlightening from a physical point of view.

$$V(H) \supset \sqrt{2H^\dagger H} = v + h$$

$$(v + h)^{2k+1} = \left( \sqrt{2H^\dagger H} \right)^{2k+1}$$

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Field redefinition: 
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$$V(H) \propto \frac{1}{4}(v+h)^2 + \frac{1}{4v}(v+h)^3 + \frac{1}{16v^2}(v+h)^4$$

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$$\begin{aligned} &= \left[ \frac{1}{2}(v+h) + \frac{1}{4v}(v+h)^2 \right]^2 \\ &= (v_1 + h_1)^2 = 2H_1^\dagger H_1 \end{aligned}$$

One can convert the SMEFT Lagrangian to HEFT form using Eq. (2.11) to switch from Cartesian and polar coordinates. One can attempt to convert from HEFT to SMEFT form using

$$\frac{\phi}{(\phi \cdot \phi)^{1/2}} = n \quad (2.30)$$

with  $(\phi \cdot \phi)^{1/2}$  some function of  $h$ . This substitution gives a Lagrangian  $L(\phi)$  that need not be analytic in  $\phi$ . However, if there is an  $O(4)$  fixed point, then there is a suitable change of variables such that the resulting Lagrangian is analytic in  $\phi$ .

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} \phi^1 + i\phi^2 \\ \phi^3 + i\phi^4 \end{pmatrix}$$

$$\mathcal{L}_{\text{Kinetic}} = \frac{1}{2} g_{ij}(\phi) (\partial_\mu \phi^i) (\partial^\mu \phi^j)$$



$$ds^2 \equiv g_{ij}(\phi) d\phi^i d\phi^j$$

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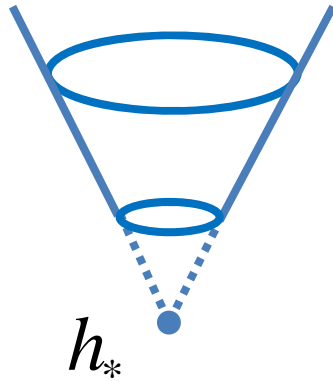
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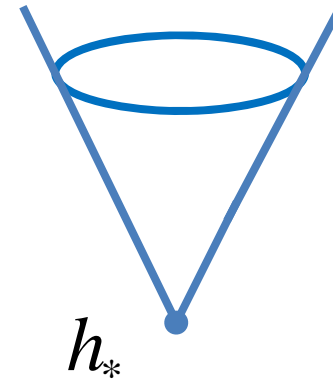
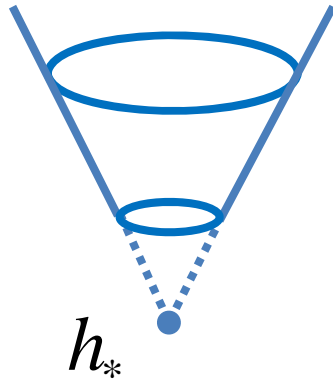
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$$\exists h_* \text{ such that } F(h_*) = 0$$



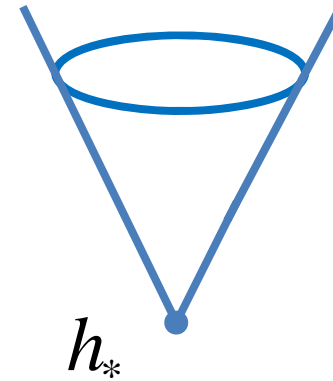
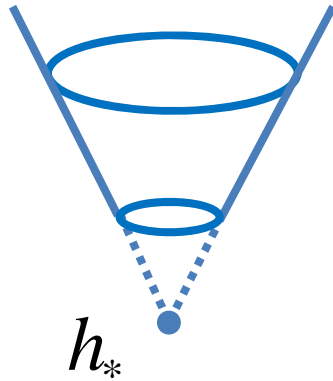
AJM (arXiv: 1605.03602)

$\exists h_*$  such that  $F(h_*) = 0$



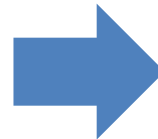
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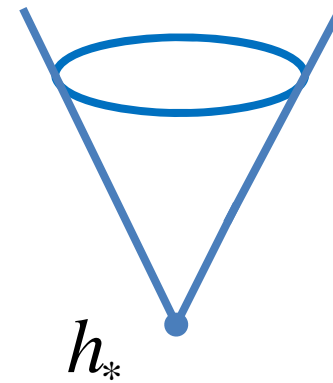
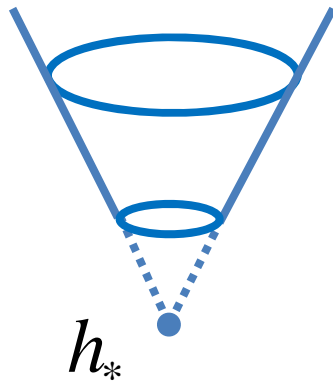
us (arXiv: 2008.08597)

$$F = 0$$

$$R, \nabla^2 R, \nabla^4 R, \dots < \infty$$

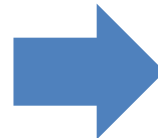
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# UV theories that will generate HEFT?



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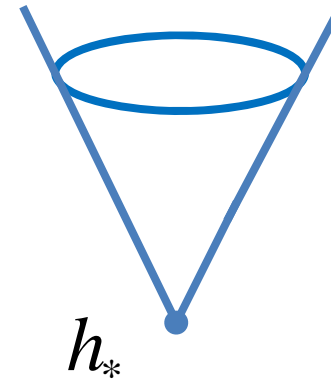
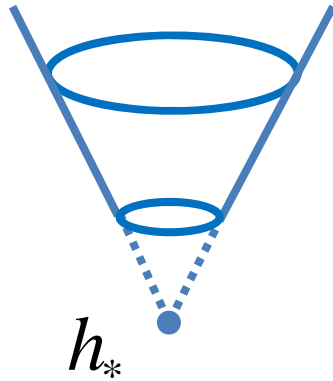
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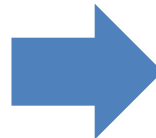
# UV theories that will generate HEFT?

- Extra electroweak symmetry breaking



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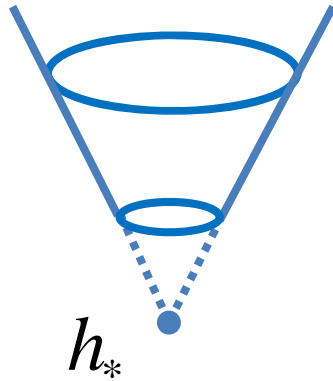
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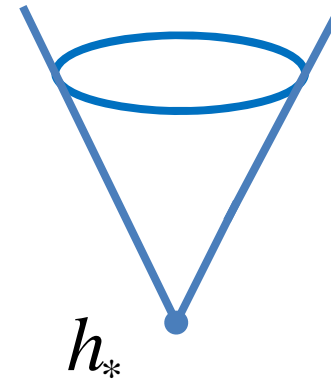
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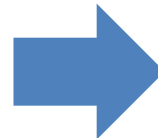


- Mass fully from electroweak symmetry breaking



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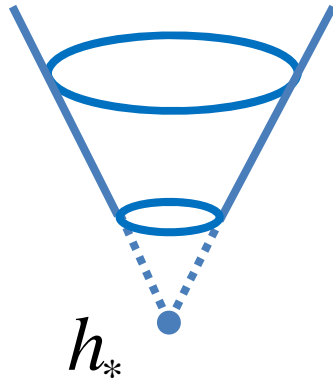
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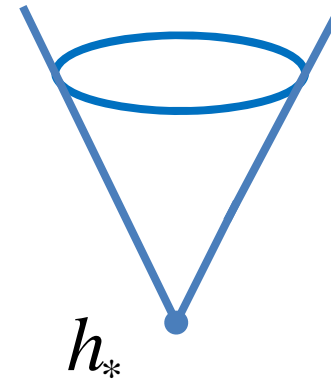
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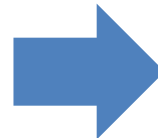
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Leading Order  
Criterion

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us (arXiv: 2008.08597)

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## Matching to **all-order in fields**

Coleman-Weinberg:  $\mathcal{L}_{\text{UV}} \supset -\frac{1}{2}\Phi\left[\partial^2 + M^2 + U(\phi)\right]\Phi$

$$i \ln \det(\partial^2 + M^2 + U) = i \text{Tr} \ln(\partial^2 + M^2 + U)$$

$$= i \int d^4x \int \frac{d^4p}{(2\pi)^4} \text{tr} \left[ \ln(-p^2 + M^2 + U) \right]$$

$$= \int d^4x \frac{1}{16\pi^2} \text{tr} \left[ \frac{1}{2} (M^2 + U)^2 \left( \ln \frac{\mu^2}{M^2 + U} + \frac{3}{2} \right) \right]$$

## Matching to **all-order in fields**

Coleman-Weinberg:  $\mathcal{L}_{\text{UV}} \supset -\frac{1}{2}\Phi\left[\partial^2 + M^2 + U(\phi)\right]\Phi$

$$\begin{aligned}i \ln \det(\partial^2 + M^2 + U) &= i \text{Tr} \ln(\partial^2 + M^2 + U) = i \int d^4x \int \frac{d^4p}{(2\pi)^4} \text{tr} \ln\left[-(p + i\partial)^2 + M^2 + U\right] \\ &= i \int d^4x \int \frac{d^4p}{(2\pi)^4} \text{tr}\left[\ln(-p^2 + M^2 + U)\right] \\ &= \int d^4x \frac{1}{16\pi^2} \text{tr}\left[\frac{1}{2}(M^2 + U)^2 \left(\ln \frac{\mu^2}{M^2 + U} + \frac{3}{2}\right)\right]\end{aligned}$$

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 \end{aligned}$$

If  $[U, \partial_\mu U] \neq 0$ : see App. D in 2008.08597



## Example: A heavy Singlet

$$\mathcal{L}_{\text{UV}} = \mathcal{L}_{\text{SM}} - \frac{1}{2} S \left( \partial^2 + M^2 + \kappa |H|^2 \right) S \quad , \quad m_S^2 = M^2 + \frac{1}{2} \kappa v^2$$

$M^2 > 0$  , SMEFT       $M^2 = 0$  , HEFT

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$$\mathcal{L}_{\text{EFT}} = \mathcal{L}_{\text{SM}} + \frac{1}{16\pi^2} \left[ \frac{1}{4} \left( M^2 + \kappa |H|^2 \right)^2 \left( \ln \frac{\mu^2}{M^2 + \kappa |H|^2} + \frac{3}{2} \right) + \frac{1}{24} \frac{\kappa^2}{M^2 + \kappa |H|^2} \left( \partial |H|^2 \right)^2 \right]$$

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$$R = \frac{1}{16\pi^2} \left[ \frac{1}{2K^4} \frac{\kappa^2 M^2}{\left( M^2 + \kappa |H|^2 \right)^2} + \frac{1}{2K^2} \frac{\kappa^2}{M^2 + \kappa |H|^2} \right] \quad , \quad K^2 \equiv 1 + \frac{1}{96\pi^2} \frac{\kappa^2 |H|^2}{M^2 + \kappa |H|^2}$$

# Summary

- A geometric criterion to tell a HEFT Lagrangian from SMEFT
- An understanding of what UV theories would generate HEFT
- A functional method to match to all order in *fields* (not *derivatives* yet)
- A few UV examples show that the leading order criterion works